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• Non-magnetic structure considered $\mu_r = 1$

• Different permittivity tensors $\overline{\overline{\epsilon}}_n$ for each layer $n = 1, \cdots, N$ • Uniaxial dielectric (glass-based) or conductive (graphite-based) media • Large-scale hypothesis: each layer is homogeneous $\bar{\epsilon}_n = \operatorname{diag}\left[\varepsilon_{n;11}; \varepsilon_{n;22}; \varepsilon_{n;22}\right]$ in the material framework



delaminations (thin, air-type gap)

• Uniaxial defects: 3-D (voids, fluid-filled cavities, cracks, etc.), or

	•
<u>E</u> _N	

Approximate solution of the direct problem

First-order solution

• Asymptotic formula of the field scattered by an inclusion $\overline{\overline{\epsilon}}_i = \overline{I}\epsilon_i$ at the \mathbf{r}_m position

 $\mathbf{E}^{\mathrm{sca}}(\mathbf{r}) = \mathrm{i}\omega\mu_0 \bar{\bar{G}}(\mathbf{r}, \mathbf{r}_m) \cdot \bar{\bar{\varrho}} \cdot \mathbf{E}^{\mathrm{inc}}(\mathbf{r}_m)$

• Polarization tensor $\bar{\bar{\varrho}} = -i\omega\epsilon_0 V \text{diag}\left[\alpha_l, \alpha_t, \alpha_t\right]$ • $\alpha_l = \varepsilon_{n;11} \frac{\epsilon_i - \varepsilon_{n;11}}{\varepsilon_{n;11} + L_l (\epsilon_i - \varepsilon_{n;11})}$, $\alpha_t = \varepsilon_{n;22} \frac{\epsilon_i - \varepsilon_{n;22}}{\varepsilon_{n;22} + L_t (\epsilon_i - \varepsilon_{n;22})}$ • For a cubical inclusion of side a and volume $V = a^3$: $L_l = c \arctan(c / \sqrt{1 + 2c})$ et $L_l + 2L_t = 1$, $c = \varepsilon_{n;11} / \varepsilon_{n;22}$ are found in W. S. Weiglhofer and A. Lakhtakia, Int. J. Infrared Millim. Waves **17** 1365-1376 (1996)

Spectral response

• Numerical solution of the state equation $\frac{d}{dz}\psi(z) = \overline{A}_n \cdot \psi(z) + \mathbf{f}(z)$ where $\psi(z)$ is the tangential field vector and $\mathbf{f}(z)$ is the source vector

$$\boldsymbol{\psi}\left(z\right) = \begin{bmatrix} k_{x}\tilde{H}_{x}\left(z\right) + k_{y}\tilde{H}_{y}\left(z\right) \\ k_{y}\tilde{H}_{x}\left(z\right) - k_{x}\tilde{H}_{y}\left(z\right) \\ k_{x}\tilde{E}_{x}\left(z\right) + k_{y}\tilde{E}_{y}\left(z\right) \\ k_{y}\tilde{E}_{x}\left(z\right) - k_{x}\tilde{E}_{y}\left(z\right) \end{bmatrix} \qquad \qquad \mathbf{f}\left(z\right) = \begin{bmatrix} -k_{y}\tilde{J}_{x}\left(z\right) + k_{x}\tilde{J}_{y}\left(z\right) \\ k_{y}\tilde{J}_{y}\left(z\right) + k_{x}\tilde{J}_{x}\left(z\right) \\ \frac{k_{x}^{2} + k_{y}^{2}}{\omega\epsilon_{zz}}\tilde{J}_{z}\left(z\right) \\ 0 \end{bmatrix} = \begin{bmatrix} -k_{y}\tilde{J}_{x}\left(z\right) + k_{x}\tilde{J}_{y}\left(z\right) \\ k_{y}\tilde{J}_{y}\left(z\right) + k_{x}\tilde{J}_{x}\left(z\right) \\ \frac{k_{x}^{2} + k_{y}^{2}}{\omega\epsilon_{zz}}\tilde{J}_{z}\left(z\right) \\ 0 \end{bmatrix}$$

• ODE solution via eigenvalue decomposition of the so-called propagator matrix \bar{A}_n



Numerical results - glass-fiber case

- $f = 6 \,\mathrm{GHz}$
- $\bar{\bar{\epsilon}}_1 = \epsilon_0 \text{diag} \left[5.46 + i2.29, 5.21 + i2.08, 5.21 + i2.08 \right]$
- $\theta_1 = 0^\circ$ with respect to the optical axis x
- Receivers-emitters: array made of 13 electric dipoles, xand y orientation
- Inclusion: air cube of side $0.1\lambda_0$ at $\mathbf{r}_m = (0, 0, -\lambda_0)$

Observed scattered electric field. Emitter is the central dipole with x and y orientation (rows)



$$\frac{d}{dz}\boldsymbol{\psi}\left(z\right) = \bar{\bar{U}}_{n}\cdot\bar{\bar{A}}_{n}\cdot\bar{\bar{U}}_{n}^{-1}\cdot\boldsymbol{\psi}\left(z\right) + \mathbf{f}\left(z\right)$$

Spatial response

• The goal is to compute the I-FT of fast-oscillating functions in the $k_x - k_y$ plane

$$G(x,y) = \frac{1}{4\pi^2} \iint_{-\infty}^{\infty} \tilde{G}_0(k_x,k_y) \mathrm{e}^{\mathrm{i}(k_x x + k_y y)} \,\mathrm{d}k_x \,\mathrm{d}k_y$$

• Interpolation of the stable part through the Pad_n degree of Padua points with scaled Chebyshev polynomials

$$\mathcal{L}_{\mathrm{Pad}_{n}}\tilde{G}_{0}\left(k_{x},k_{y}\right) = \sum_{k=0}^{\mathrm{Pad}_{n}}\sum_{j=0}^{k}c_{j,k-j}\hat{T}_{j}\left(k_{x}\right)\hat{T}_{k-j}\left(k_{y}\right) - \frac{1}{2}c_{\mathrm{Pad}_{n},0}\hat{T}_{\mathrm{Pad}_{n}}\left(k_{x}\right)\hat{T}_{0}\left(k_{y}\right)$$

with the coefficients $c_{j,k-j}$ given by M. Caliari *et al.*, Num. Alg. **56** 45-60 (2011) • FT of Chebyshev polynomials given by V. Domínguez et al., IMA J. Num. Anal. **31** 1253-1280 (2011)

$$\int_{-1}^{1} \hat{T}_{j}(k_{x}) e^{-ik_{x}x} dk_{x} = \sqrt{2} \int_{-1}^{1} \cos\left[j\cos^{-1}(k_{x})\right] e^{-ik_{x}x} dk_{x}$$

Figure: Dipole along the y axis: E_x (— MoM, -•- asymptotic), E_y (— MoM, -•- asymptotic), E_z (— MoM, -•- asymptotic).







Conclusions and perspectives

- Numerical integration method based on the Padua-Domínguez interpolation-integration technique is proposed to efficiently compute integration of fast oscillating functions
- The approach is validated by comparisons with a Method of Moments (MoM) in Y. Zhong et al., IEEE Trans. Antennas Prop., submitted (2014)
- Numerical examples with multi-layer structures made of glass & carbon fiber materials
- Uniaxial defects with various orientations
- Validation by comparisons with experimental data