



HAL
open science

Low temperature plasma enhanced deposition of GaP films on Si substrate

Alexander Gudovskikh, Ivan Morozov, Alexander Uvarov, Dmitriy Kudryashov, Ekaterina Nikitina, Anton Bukatin, Vladimir Nevedomskiy, Jean-Paul Kleider

► **To cite this version:**

Alexander Gudovskikh, Ivan Morozov, Alexander Uvarov, Dmitriy Kudryashov, Ekaterina Nikitina, et al.. Low temperature plasma enhanced deposition of GaP films on Si substrate. *Journal of Vacuum Science & Technology A*, 2018, 36 (2), 10.1116/1.4999409 . hal-01777504

HAL Id: hal-01777504

<https://centralesupelec.hal.science/hal-01777504v1>

Submitted on 25 Mar 2020

HAL is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

Low temperature plasma enhanced deposition of GaP films on Si substrate

Alexander S. Gudovskikh, Ivan A. Morozov, Alexander V. Uvarov, Dmitriy A. Kudryashov, Ekaterina V. Nikitina, Anton S. Bukatin, Vladimir N. Nevedomskiy, and Jean-Paul Kleider

Citation: *Journal of Vacuum Science & Technology A* **36**, 021302 (2018); doi: 10.1116/1.4999409

View online: <https://doi.org/10.1116/1.4999409>

View Table of Contents: <http://avs.scitation.org/toc/jva/36/2>

Published by the [American Vacuum Society](#)

Articles you may be interested in

[Synergistic interactions of H₂ and N₂ with molten gallium in the presence of plasma](#)

Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films **36**, 021303 (2018);
10.1116/1.5004540

[Characterization of particle generated during plasma-enhanced chemical vapor deposition on amorphous carbon layer using particle beam mass spectrometer](#)

Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films **36**, 021506 (2018);
10.1116/1.5000087

[Nucleation mechanism during WS₂ plasma enhanced atomic layer deposition on amorphous Al₂O₃ and sapphire substrates](#)

Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films **36**, 01A105 (2018);
10.1116/1.5003361

[Sub-damage-threshold plasma etching and profile tailoring of Si through laser-stimulated thermal desorption](#)


Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films **36**, 021301 (2018);
10.1116/1.4991586

[Plasma enhanced atomic layer deposition of aluminum sulfide thin films](#)

Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films **36**, 01A113 (2018);
10.1116/1.5003339

[Low temperature plasma enhanced deposition approach for fabrication of microcrystalline GaP/Si superlattice](#)


Journal of Vacuum Science & Technology A: Vacuum, Surfaces, and Films **36**, 02D408 (2018);
10.1116/1.5018259



Instruments for Advanced Science

Contact Hiden Analytical for further details:
W www.HidenAnalytical.com
E info@hiden.co.uk

CLICK TO VIEW our product catalogue



Gas Analysis

- dynamic measurement of reaction gas streams
- catalysis and thermal analysis
- molecular beam studies
- dissolved species probes
- fermentation, environmental and ecological studies



Surface Science

- UHV TPD
- SIMS
- end point detection in ion beam etch
- elemental imaging - surface mapping



Plasma Diagnostics

- plasma source characterization
- etch and deposition process reaction kinetic studies
- analysis of neutral and radical species



Vacuum Analysis

- partial pressure measurement and control of process gases
- reactive sputter process control
- vacuum diagnostics
- vacuum coating process monitoring

Low temperature plasma enhanced deposition of GaP films on Si substrate

Alexander S. Gudovskikh,^{a)} Ivan A. Morozov, Alexander V. Uvarov, Dmitriy A. Kudryashov, Ekaterina V. Nikitina, and Anton S. Bukatin

Nanotechnology Centre, St. Petersburg Academic University RAS, Hlopina Str. 8/3, St.-Petersburg 194021, Russia

Vladimir N. Nevedomskiy

A. F. Ioffe Physical-Technical Institute RAS, Polytechnicheskaya Str. 26, St.-Petersburg 194021, Russia

Jean-Paul Kleider

GeePs, Group of Electrical Engineering-Paris, CNRS, CentraleSupélec, Univ. Paris-Sud, Université Paris-Saclay, Sorbonne Universités, UPMC Univ Paris 06, 11 rue Joliot-Curie, 91192 Gif-sur-Yvette Cedex, France

(Received 8 August 2017; accepted 13 November 2017; published 8 December 2017)

Amorphous and microcrystalline GaP films were deposited on Si substrates by time modulated plasma enhanced deposition from trimethylgallium and phosphine using constant hydrogen plasma at a temperature of 250–380 °C. Amorphous GaP films obtained at constant low radio-frequency (RF) power (20 W) mode exhibit the broad feature at 350–360 cm⁻¹ and a shoulder at 370–390 cm⁻¹ in Raman spectra. Amorphous films have smooth surface with root-mean-square (RMS) roughness decreasing from 0.9 to 0.2 nm with increasing deposition temperature from 250 to 380 °C. Small amounts of 3–5 nm GaP nanocrystal inclusions in the amorphous matrix are formed at 380 °C. An increase of RF power to 100 W during Ga and P deposition steps leads to the formation of a GaP crystalline phase as confirmed by transmission electron microscopy. Two peaks in the Raman spectra at 365 and 402 cm⁻¹, which correspond to GaP TO-LO duplet, were observed in this case. However, the microcrystalline GaP layers have rough surface with RMS roughness of 6 nm. A critical role of the hydrogen plasma for GaP growth mechanism is demonstrated. Indeed an effect of the phosphorous accumulation in the deposition chamber during phosphine plasma assisted decomposition step is shown. The hydrogen plasma treatment after P deposition step allows one to avoid an influence of excess phosphorous to the growth process providing a way to obtain ALD like growth of GaP with good thickness uniformity. *Published by the AVS.* <https://doi.org/10.1116/1.4999409>

I. INTRODUCTION

Heterostructures with GaP layer grown on Si substrates are of great interest for optoelectronic devices integrated with silicon, in particular, for solar cells. GaP has the smallest lattice mismatch beyond III-V binary compounds (less 0.4%) being the best candidate for nucleation (buffer) layer providing conditions for low defect density at the interface with the Si substrate.¹ GaP has a gap of 2.26 eV fitting requirements for a wide gap emitter of GaP/Si heterojunction solar cells, which according to simulations could have an advantage of more than 1% of efficiency compared to passivated emitter and rear cells made of p-type Si wafers,² which are still of interest, in particular, for low orbit space applications due to better radiation hardness.³ Recently, promising results have been obtained using n-GaP window for Si n-p homojunctions.⁴ However, GaP layers are commonly grown by epitaxial techniques such as molecular beam epitaxy (MBE)^{5,6} and metal organic vapor-phase epitaxy (MOVPE),^{7–9} which require relatively high temperatures (500–800 °C). Until now, the photovoltaic performance of GaP/Si heterojunctions obtained by epitaxy^{4,10,11} is far from that of high efficiency Si solar cells. The main problem, especially for structures with anisotype heterojunctions (n-GaP/p-Si), can be related to the high growth temperature, which affects the GaP/Si interface quality.¹² The use of high

temperatures could lead to interdiffusion of group-III and -V atoms into Si and opposite, which act as dopants affecting the electrical properties of heterojunctions, and it can also promote the fast diffusion of species that can degrade the carrier lifetime in c-Si.¹³

In this paper, we explore a low temperature (less than 400 °C) technological approach for the growth of III-V compounds on Si substrates using plasma-enhanced atomic layer deposition (PE-ALD) approach. The PECVD related low-temperatures technology could provide a lot of advantages for PV mass production such as large area, high throughput, low price of the equipment, and the process. Low temperature ALD is becoming widely used in photovoltaics for deposition of dielectric antireflection and passivation layers (commonly Al₂O₃) on the silicon surface¹⁴ while only few reports on plasma deposition of III-V semiconductors at low temperature could be found. However, the first report on plasma enhanced growth at 300 and 350 °C of polycrystalline GaN and amorphous GaP films,¹⁵ respectively, is dated by 1978. Recently, the GaN layers containing small crystallites dispersed in an amorphous matrix were grown on Si using PE-ALD in the temperature range of 185–385 °C.¹⁶ Our preliminary results on plasma enhanced deposition of amorphous GaP at 350 °C have demonstrated promising passivation properties of GaP/Si heterostructures when the phosphorus and gallium atom source flows were alternatively changed.¹⁷ However, the obtained deposition rate of

^{a)}Electronic mail: gudovskikh@spbau.ru

TABLE I. Parameters of the process steps: time, gas flows, and RF power.

Step	Constant RF power (20 W) mode				Pulsed RF power mode				
	t (s)	Gas flow (sccm)			t (s)	Gas flow (sccm)			RF (W)
		PH ₃	H ₂	TMG/H ₂		PH ₃	H ₂	TMG/H ₂	
P step	3–5	50	50	—	2–3	50	50	—	100
H plasma treatment	—	—	—	—	0–5	—	100	—	100
Purge	10	—	100	—	7	—	100	—	20
Ga step	5–10	—	30	70	3–5	—	30	70	100
H plasma treatment	—	—	—	—	0–2	—	100	—	100
Purge	10	—	100	—	7–9	—	100	—	20

11.6 nm/cycle means that the ALD growth mode was not reached. Moreover, Ar plasma, which was used for plasma ignition and discharge, produces radiation defects in Si mitigating the advantage of low-temperature approach. Thus, PE-ALD like technique with hydrogen plasma is used in this study.

II. EXPERIMENT

The time modulated plasma enhanced deposition (PE-ALD like approach) was realized using a standard PECVD setup with capacitively coupled plasma (Oxford Plasmalab 100) supplied with phosphine (PH₃) and trimethylgallium (TMG) lines for sources of phosphorus and gallium, respectively. Hydrogen was used as a gas carrier for TMG. The hydrogen pressure was 2 bar while the temperature of the TMG cylinder was kept at 3 °C. Silicon substrates with (001) orientation were used for GaP deposition. At the initial stage, the native oxide was removed from the Si surface by HF-dip (10%) immediately before loading to the PECVD chamber (via a load-lock system). The substrates were heated to 250–380 °C during 20–30 min in the vacuum chamber. The total time between HF dip and start of GaP layer deposition process does not exceed 40 min. The ultimate pressure was provided by BOC Edwards iH600 dry vacuum pump, and was estimated to be less 0.1 mTorr (below measurement limit of Pirani gauge).

In order to avoid the Ar plasma influence to Si substrate, the process was performed with continuous H₂ plasma discharge during deposition and purge cycles. No evacuation cycles were used because of the difficult plasma ignition of hydrogen rich gas mixture. Due to specific features of our PECVD setup, the purge step should be at least 7 s for full gas exchange and the minimal time of any step is 2 s to provide stable program operation. Total gas flow and gas pressure were always kept constant and were equal to 100 sccm and 350 mTorr, respectively. Stable gas dynamic conditions with minimal throttle valve adjustment were reached.

First, the deposition of a series of GaP layers was carried out with temperature variation from 250 to 380 °C and at a power of 20 W (corresponds to a power density of 0.03 W/cm²) which is the minimal value required for stable plasma RF (13.56 MHz) discharge. The parameters of each step are presented in Table I. Then, the deposition with pulsed RF power increase to 100 W (0.14 W/cm²) during the precursor

decomposition and short treatment in hydrogen plasma was performed at the temperature of 380 °C. The total number of cycles for each process was equal to 200.

The layer morphology and thickness were determined using scanning electron microscopy (SEM) (SUPRA 25 Zeiss) and atomic force microscopy (AFM) (BioScope Catalyst Bruker). The structural properties were studied by Raman spectroscopy using Horiba Jobin Yvon LABRAM HR 800 spectrometer with a laser excitation at 532 nm and by transmission electron microscopy (TEM) (JEOL JEM 2100F) with 200 kV acceleration voltages. Optical emission spectroscopy (OES) was used for *in situ* gas content analysis using Solar Laser Systems S100 spectrometer.

III. RESULTS AND DISCUSSION

A. Constant RF power mode

According to x-ray fluorescence measurements, the deposited films content Ga and P atoms while no impurities could be detected because of low sensitivity of this technique to thin films. SEM [Fig. 1(a)] and TEM [Fig. 1(b)] demonstrate that GaP films deposited with a constant RF power of 20 W are homogeneous with a smooth surface and sharp GaP/Si interface. The selected area electron diffraction (SAED) pattern of the GaP film obtained by TEM [inset of Fig. 1(b)] exhibits intensive dots corresponding to diffraction from Si substrate and diffused ring (halo) at (111) plan indicating a dominant amorphous phase of GaP.

The growth per cycle for the GaP layers grown at 350 °C and above was about 0.25 ± 0.02 nm/cycle, indicating the

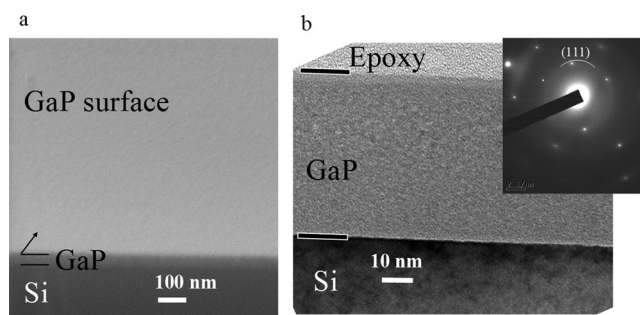


Fig. 1. SEM 20° tilt (a) and TEM bright field (b) cross section images of the GaP film grown on Si at 380 °C with constant RF power (20 W) mode. SAED pattern by TEM is presented in the inset.

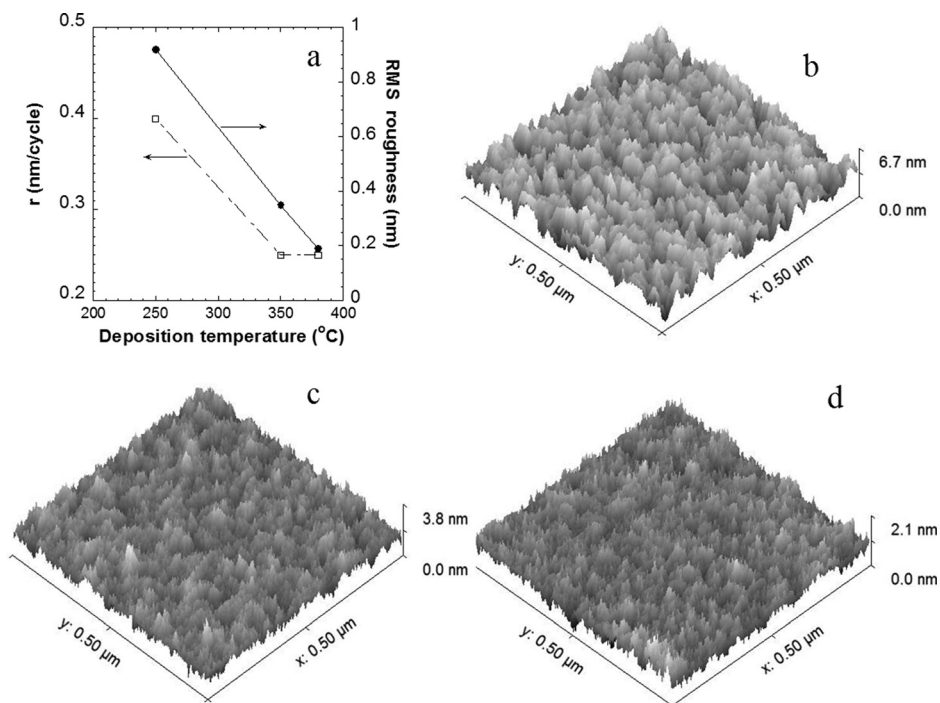


Fig. 2. Temperature dependence of the deposition rate (r) and RMS roughness (a) and surface topology by AFM for GaP deposited at 250°C (b), 350°C (c), and 380°C (d).

growth of approximately one GaP monolayer (0.275 nm) per cycle [Fig. 2(a)]. Thus, one monolayer per cycle growth mode was achieved using hydrogen plasma in contrary to results obtained with argon plasma (1–1.6 nm/cycle for 350°C). For 250°C, the growth per cycle was about 0.4 nm/cycle. The surface topography measured by AFM strongly depends on the growth temperature (Fig. 2). The root-mean-square (RMS) roughness decreased from 0.92 nm for 250°C to 0.19 nm for 380°C.

The Raman spectra presented in Fig. 3 for GaP layers deposited on Si substrates at different temperatures exhibit the broad feature at 350–360 cm^{-1} , which corresponds to GaP

TO(X)-TO(L) peak and a shoulder at 370–390 cm^{-1} . Such behavior of Raman spectra indicates an amorphous structure of the GaP films.¹⁸ An amorphous structure of PECVD GaP films grown at 350°C was also reported in Ref. 15. We should note that the shoulder at 370–390 cm^{-1} becomes more pronounced when the deposition temperature is increased, and it shifts toward the GaP LO peak position (402 cm^{-1}), indicating better structural properties. Indeed, a formation of GaP nanocrystals was observed by HRTEM in the layers deposited at 380°C. Low fraction of 2–3 nm size ordered inclusions embedded in amorphous GaP matrix could be selected in the HRTEM image presented in Fig. 4.

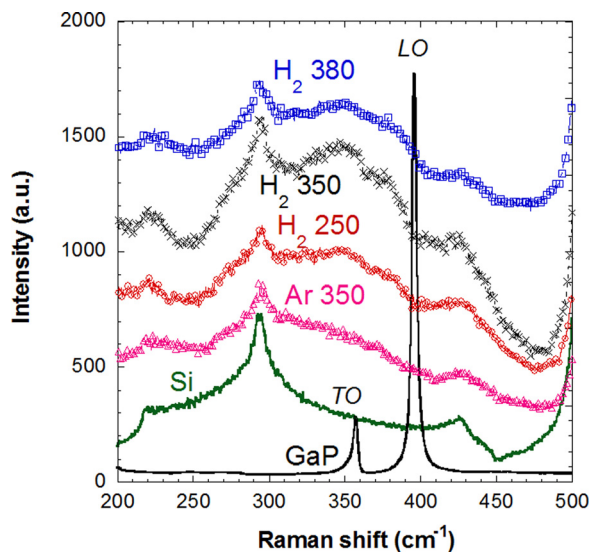


Fig. 3. (Color online) Raman spectra of GaP deposited on Si with constant RF power at different temperatures. Spectra of Si and GaP substrates are also presented.

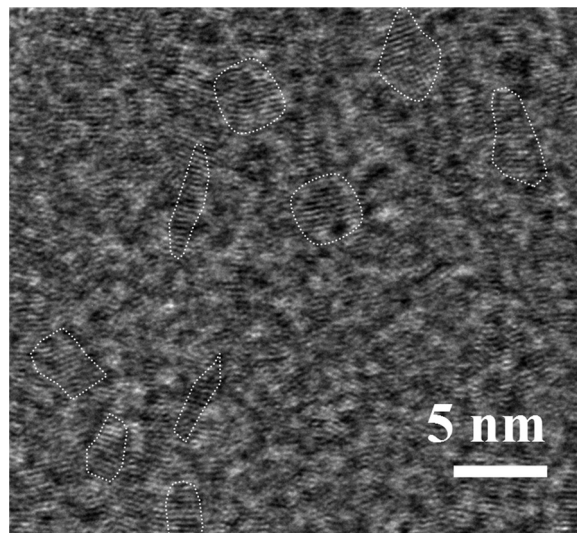


Fig. 4. HRTEM image of GaP grown at 380°C with constant RF power (20W) mode. Several examples of GaP ordered areas are marked by dot contours.

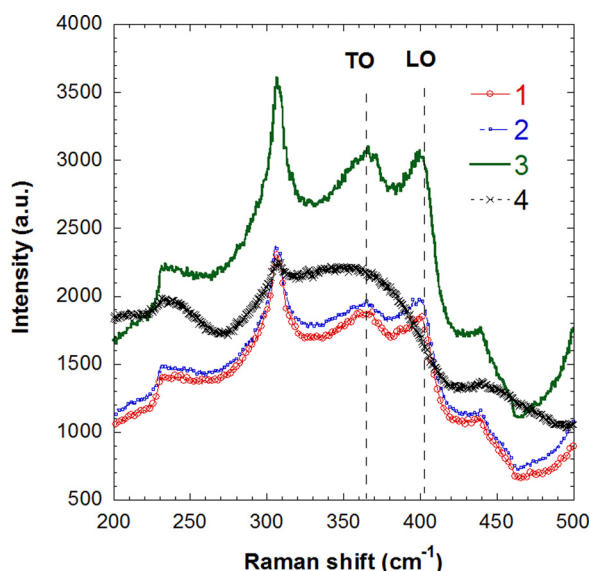


Fig. 5. (Color online) Raman spectra of GaP deposited on Si with pulsed RF power (100 W) during Ga and P deposition steps: without H₂ plasma treatment (1), with H₂ plasma (100 W) treatment step after Ga and P deposition steps (2), with H₂ plasma (100 W) treatment step after P deposition step only (3). GaP deposited at low RF power (20 W) during Ga and P deposition steps but with H₂ plasma (100 W) treatment step after Ga and P deposition step (4).

Obviously higher deposition temperature is required to reach a successful surface migration of the deposited atoms for formation of a GaP crystalline phase. However, the deposition temperature of 380 °C approaches the practical limit of our setup (400 °C) being also close to a temperature limit for the most commonly used plasma-enhanced equipment. Moreover, further temperature increase is undesirable because higher temperature could affect the electronic properties of the GaP/Si interface and Si wafers due to interdiffusion.

B. Pulsed RF power mode

To enhance the surface migration of adsorbed atoms a pulsed increase of the RF power to 100 W during Ga and P deposition steps was proposed. Also the possibility to enhance the surface migration by an additional short hydrogen treatment step at increased RF power introduced immediately after Ga and P deposition steps was explored. The RF

power during purge steps was decreased to the minimum value (20 W), which enables to hold the glow discharge and to minimize the influence of hydrogen plasma to Si. The Raman spectra for a GaP film deposited with pulsed RF power mode are presented in Fig. 5. Two clear peaks at 365 and 402 cm⁻¹ could be observed in the Raman spectra for the layers grown with increased RF power during Ga and P deposition steps. The peak positions correspond to the GaP TO-LO duplet¹⁹ revealing the formation of a GaP crystalline phase. The study of structural properties by TEM demonstrates that GaP films are microcrystalline (Fig. 6). The selected area electron diffraction (SAED) pattern of the GaP film obtained by TEM confirms the presence of a GaP crystalline phase in the layer [Fig. 6(b)]. Intensive dots in SAED pattern correspond to diffraction from Si substrate, while the sharp rings correspond to (111), (220), and (311) reflections of disoriented GaP crystals. However, high intensive areas within the diffraction rings indicate a preferred orientation of GaP crystals. A significant crystalline fraction in GaP films grown with enhanced RF power is observed in the HRTEM image presented in Fig. 6(c) where disoriented GaP grains of few nanometers size could be selected. The formation of grains leads to rough surface of the GaP film deposited at the pulsed RF power mode as illustrated by SEM and AFM in Fig. 7. The RMS roughness is within the range of 5.87–5.92 nm.

When only the short hydrogen treatment step with increased RF power (100 W) was introduced immediately after low power (20 W) Ga and P deposition steps the deposited films were amorphous as indicated by Raman spectra (Fig. 5) with smooth surface (RMS roughness < 1 nm). However, the pulsed H₂ plasma treatment step strongly affects the growth per cycle and thickness uniformity. Indeed, introduction of H₂ plasma treatment after Ga deposition step leads to the increase of growth per cycle with deposition step time (Fig. 8), which significantly surpasses one GaP monolayer per cycle. The thickness nonuniformity increases in this case up to 50% for 3 in. wafer. On the contrary, introduction of H₂ plasma treatment after P deposition step leads to the saturation of the deposition rate with the rise of the deposition step time as for Ga as well as for P (Fig. 8). The deposition rate saturates at approximately one monolayer of GaP per cycle level. The thickness

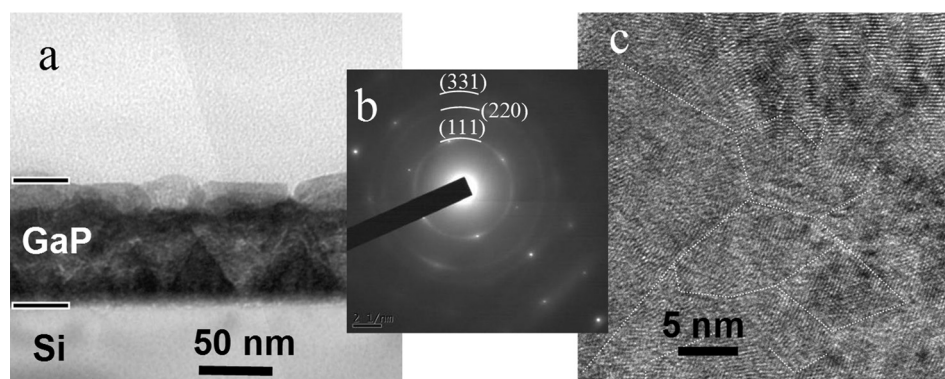


Fig. 6. TEM cross section view (bright field contrast) (a), SAED pattern (b) and HRTEM of GaP grown on Si with pulsed RF power (100 W) during deposition steps. The boundaries between disoriented GaP sections are selected by the eye.

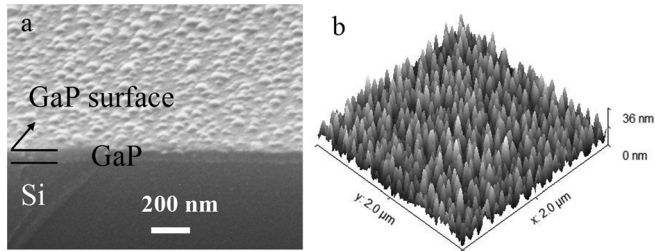


FIG. 7. SEM 20° tilt cross section image (a) and surface morphology by AFM (b) for GaP grown on Si with pulsed RF power (100 W) during deposition steps.

uniformity was extremely improved in this mode being within the range of $\pm 5\%$.

To analyze the influence of H_2 plasma treatment one should take into account two possible processes. The first one is related to gas mixture changing time because precursor residues could continue to react during the H_2 plasma treatment step. The second is related to the possible interaction between phosphorous and H_2 plasma. In particular, a red phosphorous could be condensed at the walls of the deposition chamber during PH_3 plasma decomposition.²⁰ This excess phosphorous should remain in the chamber even after a full gas mixture exchange. However, then the phosphorous is etched by H_2 plasma during Ga deposition and subsequent H_2 plasma treatment steps. In this case the formed phosphorous precursors (PH_x) would interact with Ga precursors leading to conventional PECVD growth of GaP. Therefore, the growth per cycle could exceed one monolayer per cycle with a strong spatial nonuniformity because P etching from the chamber walls provides nonhomogeneous gas content.

C. Hydrogen plasma treatment

Optical emission spectroscopy (OES) was used for *in situ* examination of the processes during hydrogen plasma treatment. The measured OES spectra of H_2 and PH_3 plasma are presented in Fig. 9. The pure hydrogen plasma has the most intensive H_α line (from Balmer series) at 656 nm caused by emission of atomic hydrogen. The short wavelength lines H_β

and H_γ at 486 and 434 nm have lower intensity. There are also spectral lines of molecular hydrogen at 600 nm range called Fulcher α band.²¹ Thus, the H_α line was used to control the hydrogen atom concentration in the plasma.

When phosphine is added to the hydrogen plasma it takes a green color due to continuous emission band of PH_3 molecules in the range of 450–650 nm. Phosphine plasma has also lines at 341, 666 and 717 nm, which correspond to emission lines of PH and PH_2 radicals. The most intense PH_2 666 nm line is of particular interest, because neither molecular nor atomic hydrogen emission lines appear at this wavelength, allowing one to extract information about the concentration of PH_2 radicals and atomic hydrogen in the PH_3/H_2 plasma.

The time dependence of the PH_2 and H_α lines during deposition process with pulsed RF power increase and H_2 plasma treatment step after Ga deposition is presented in Fig. 10(a). First, at the P deposition step (100 W) a rise of PH_3 concentration in the chamber leads to an increase of PH_2 line intensity, which is accompanied by a decrease of the intensity of the H_α line because PH_3 ionization energy (7.2 eV) is much lower compared to that of hydrogen (16.5 eV). Addition of PH_3 with a low ionization energy leads to redistribution of the electron energy in the plasma and therefore to a decrease of ionized hydrogen atoms concentration.²⁰ Studying changes in the intensity of the H_α hydrogen line is an indirect way to determine the variation of phosphine concentration. Next, during the purge step (20 W), the H_α line intensity increases due to the drop of PH_3 concentration in the chamber. Finally, at the Ga deposition step (100 W), a slight increase of the PH_2 line intensity is observed because of the phosphorous etching in H_2 plasma. The traces of phosphorous precursors during Ga deposition step obviously lead to uncontrollable GaP deposition.

To confirm this P etching effect, the phosphorous was deposited at the chamber walls by PH_3 decomposition during 30 s at 100 W. Then, the chamber was purged by H_2 during 1 min without plasma. Next, the H_2 plasma (100 W) etching was started. Time dependence of H_α and PH_2 OES lines is presented in Fig. 9(b). The PH_2 line was clearly observed in the plasma which intensity decreased during 20 s accompanied by the rise of the H_α line intensity. The observed

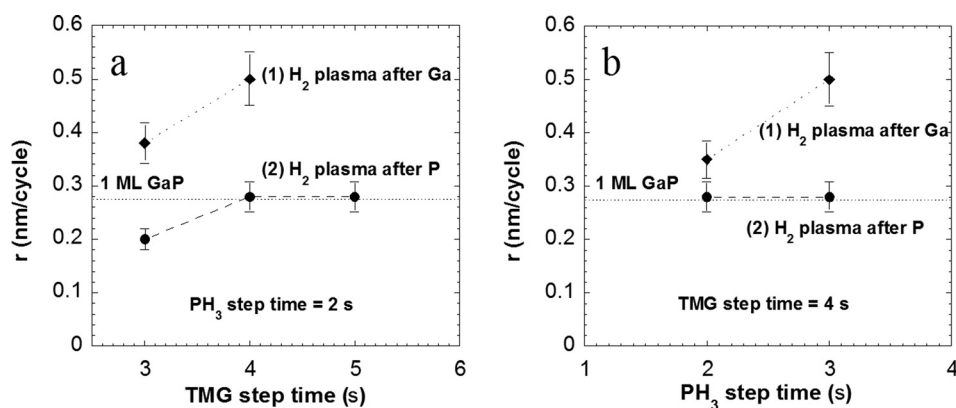


FIG. 8. Deposition rate (r) as a function of TMG (a) and PH_3 (b) deposition step time for GaP deposited with pulsed RF power (100 W) using H_2 plasma treatment after Ga deposition step (1) and H_2 plasma treatment after P deposition step. PH_3 (a) and TMG (b) deposition step times were fixed to 2 and 4 s, respectively. The level of one GaP monolayer (ML) is also indicated.

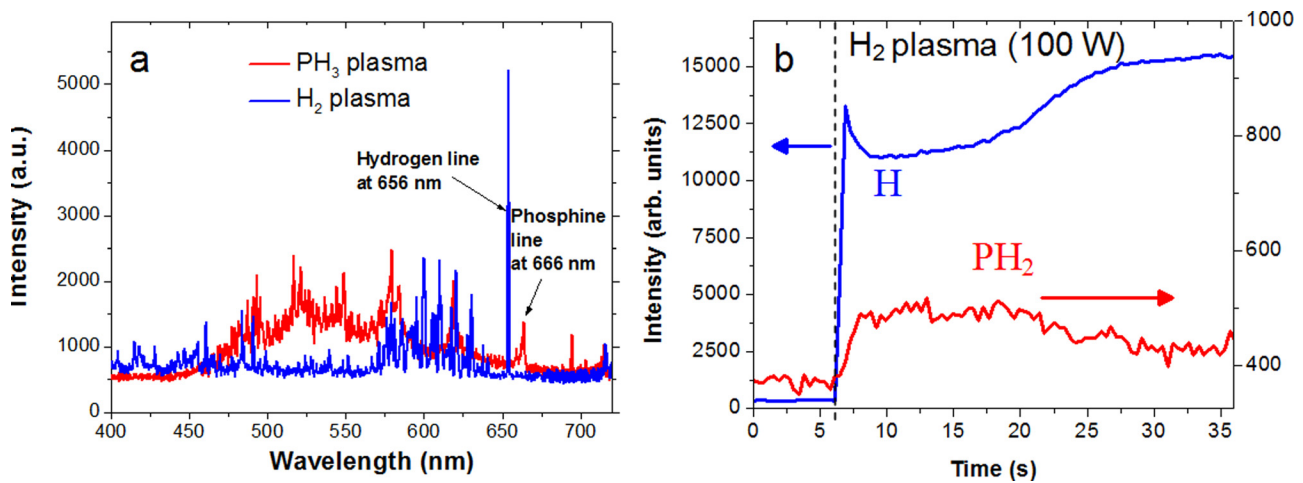


FIG. 9. (Color online) OES spectra of H₂ and PH₃ plasma (a). Time dependence of H₂ (656 nm) and PH₂ (666 nm) OES lines for H₂ plasma treatment after P deposition during 30 s (b).

behavior indicates a strong phosphorous etching effect which affects the GaP deposition process.

To avoid the negative influence of etching effect to the PE-ALD, the process was modified. Longer H₂ plasma treatment steps (100 W) were introduced after the P deposition step to etch the phosphorous in excess before Ga deposition. Moreover, no H₂ plasma treatment steps after Ga was used. The time dependence of PH₂ and H_x lines for this modified process is presented in Fig. 10(b). A decrease of PH₂ and increase of H_x lines intensity is observed during H₂ plasma treatment steps indicating a removal of excess phosphorous, while no phosphorus precursor presence (PH₂ line) is observed during the Ga deposition step. This modified pulsed RF power process allows one to reach ALD like growth of GaP films with nonuniformity less than 5% for 3 in. substrates and less than 3% for 2 in. substrates. The results of optical measurements of the uniformity are presented in Fig. 11, where the distribution of the wavelength of reflection maximum as a function of distance for 3 in. Si substrates is shown. The position of the reflection maximum, which slightly depends on distance, was used to calculate the GaP film thickness. The

thicknesses values correspond well to the one obtained by electron microscopy and are varied in narrow range indicating good film thickness uniformity.

IV. SUMMARY AND CONCLUSIONS

Plasma enhanced atomic layer deposition approach allows one to obtain amorphous or microcrystalline GaP films at temperatures below 400 °C. The transition from amorphous to microcrystalline structure could be obtained by the increase of the RF power, which increases the kinetic energy of the adsorbed atoms and their migration possibility. However, a significant increase of surface roughness appears for microcrystalline GaP films compared to the smooth surface of amorphous GaP.

A specific feature of the GaP growth mechanism was observed for higher RF power modes, namely, an accumulation of the phosphorous at the deposition chamber walls. The excess phosphorous in the chamber etched by the hydrogen plasma may interact with Ga precursors thus perturbing the ALD mode. The way to minimize this negative effect by an

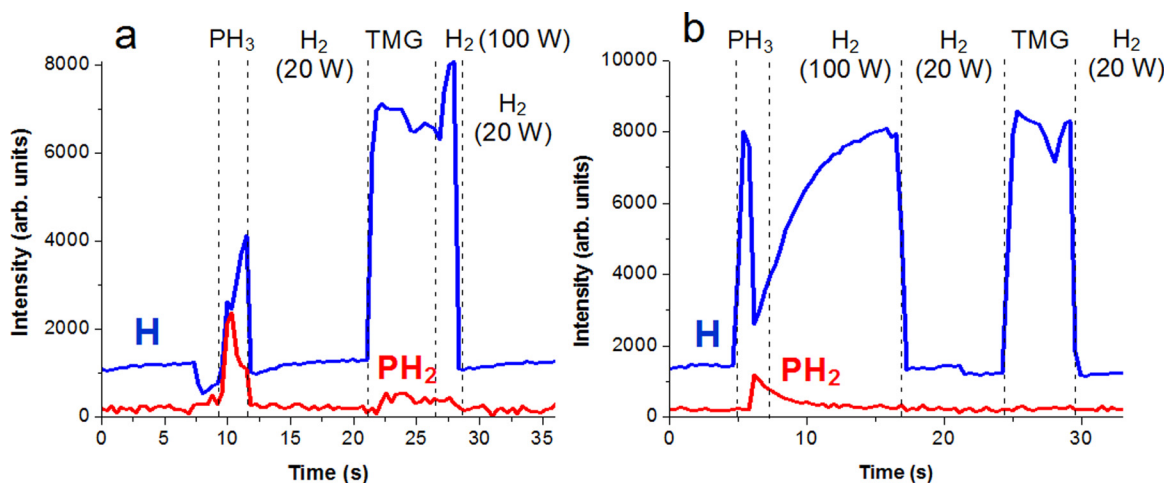


FIG. 10. (Color online) Time dependence of H_x (656 nm) and PH₂ (666 nm) OES lines for the processes with H₂ plasma treatment step after Ga deposition step (a) and after P deposition step (b).

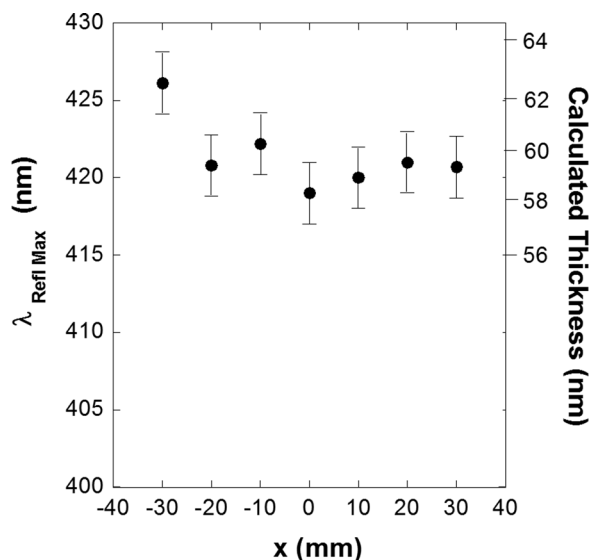


FIG. 11. Wavelength of the reflection maximum ($\lambda_{\text{Refl Max}}$) vs distance for 3 in. Si substrates with GaP film deposited with H_2 plasma treatment step after P deposition step. Calculated GaP film thickness, which corresponds to $\lambda_{\text{Refl Max}}$, is indicated on the left axis.

additional hydrogen plasma etching step after P deposition was proposed. Microcrystalline GaP films with good thickness uniformity (thickness within $\pm 5\%$ for 3 in. wafer) and saturated deposition rate (at one monolayer per cycle) were obtained. No phosphine traces in the plasma during Ga deposition step were observed in this case. Further studies will be devoted to the improvement of the structural properties and morphology of GaP films grown on Si by plasma enhanced deposition.

ACKNOWLEDGMENTS

This work was supported by the Russian Scientific Foundation under Grant No. 17-19-01482. TEM

characterizations were performed using equipment owned by the Federal Joint Research Center.

- ¹S. M. Sze, *Physics of Semiconductor Devices*, 2nd ed. (Wiley, New York, 1981).
- ²H. Wagner, T. Ohrdes, A. Dastgheib-Shirazi, B. Puthen Veetil, D. König, and P. P. Altermatt, *J. Appl. Phys.* **115**, 044508 (2014).
- ³W. C. Cooley and R. J. Janda, *Handbook of Space-Radiation Effects on Solar-Cell Power Systems* (NASA, Washington, D.C., 1963).
- ⁴M. Feifel, T. Rachow, J. Benick, and J. Ohlmann, *IEEE J. Photovoltaics* **6**, 384 (2016).
- ⁵M. Sadeghi and S. Wang, *J. Cryst. Growth* **227–228**, 279 (2001).
- ⁶S. L. Wright, H. Kroemer, and M. Inada, *J. Appl. Phys.* **55**, 2916 (1984).
- ⁷J. P. André, J. Hallais, and C. Schiller, *J. Cryst. Growth* **31**, 147 (1975).
- ⁸L. Samuelson, P. Omling, and H. G. Grimmeiss, *J. Cryst. Growth* **68**, 340 (1984).
- ⁹J. M. Olson, M. M. Al-Jassim, A. Kibbler, and K. M. Jones, *J. Cryst. Growth* **77**, 515 (1986).
- ¹⁰M. S. Sobolev, A. A. Lazarenko, E. V. Nikitina, E. V. Pirogov, A. S. Gudovskikh, and A. Yu. Egorov, *Semicond.* **49**, 559 (2015).
- ¹¹J. F. Geisz, J. M. Olson, D. J. Friedman, K. M. Jones, R. C. Reedy, and M. J. Romero, *31th IEEE PVSC* (2005), p. 695.
- ¹²A. S. Gudovskikh, K. S. Zelentsov, A. I. Baranov, D. A. Kudryashov, I. A. Morozov, E. V. Nikitina, and J.-P. Kleider, *Energy Procedia* **102**, 56 (2016).
- ¹³R. Varache, M. Darnon, M. Descazeaux, M. Martin, T. Baron, and D. Muñoz, *Energy Procedia* **77**, 493 (2015).
- ¹⁴J. A. van Delft, D. Garcia-Alonso, and W. M. M. Kessels, *Semicond. Sci. Technol.* **27**, 074002 (2012).
- ¹⁵J. C. Knights and R. A. Lujan, *J. Appl. Phys.* **49**, 1291 (1978).
- ¹⁶C. Ozgit, I. Donmez, M. Alevli, and N. Biyikli, *J. Vac. Sci. Technol., A* **30**, 01A124 (2012).
- ¹⁷I. A. Morozov, A. S. Gudovskikh, D. A. Kudryashov, E. V. Nikitina, J.-P. Kleider, A. V. Myasoedov, and V. Levitskiy, *J. Phys.: Conf. Ser.* **741**, 012088 (2016).
- ¹⁸M. Wihl, M. Cardona, and J. Tauc, *J. Non-Cryst. Solids* **8–10**, 172 (1972).
- ¹⁹S. Hayashi, *Solid State Commun.* **56**, 375 (1985).
- ²⁰G. Bruno, M. Losurdo, and P. Capezzuto, *J. Vac. Sci. Technol., A* **13**, 349 (1995).
- ²¹B. Xiao, S. Kado, S. Kajita, and D. Yamasaki, *Plasma Phys. Controlled Fusion* **46**, 653 (2004).