



**HAL**  
open science

# Proton diffusion mechanisms in the double perovskite cathode material $\text{GdBaCo}_2\text{O}_{5.5}$ : A molecular dynamics study

Fabien Briec, Guilhem Dezanneau, Marc Hayoun, Hichem Dammak

## ► To cite this version:

Fabien Briec, Guilhem Dezanneau, Marc Hayoun, Hichem Dammak. Proton diffusion mechanisms in the double perovskite cathode material  $\text{GdBaCo}_2\text{O}_{5.5}$ : A molecular dynamics study. *Solid State Ionics*, 2017, 309, pp.187 - 191. 10.1016/j.ssi.2017.07.017 . hal-01932509

**HAL Id: hal-01932509**

**<https://centralesupelec.hal.science/hal-01932509v1>**

Submitted on 23 Nov 2018

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

## Proton diffusion mechanisms in the double perovskite cathode material GdBaCo<sub>2</sub>O<sub>5.5</sub>: a molecular dynamics study

Fabien Briec<sup>a,\*</sup>, Guilhem Dezanneau<sup>a</sup>, Marc Hayoun<sup>b</sup>, Hichem Dammak<sup>a</sup>

<sup>a</sup> *Laboratoire Structures, Propriétés et Modélisation des Solides, CentraleSupélec / CNRS, Université Paris-Saclay, Grande voie des vignes, 92295 CHATENAY-MALABRY CEDEX, FRANCE*

<sup>b</sup> *Laboratoire des Solides Irradiés, École Polytechnique / CEA / CNRS, Université Paris-Saclay, 91128 PALAISEAU, France*

**Keywords:** PCFC, Proton conductor, Double perovskite, Mixed conductivity, Molecular dynamics

### Abstract

GdBaCo<sub>2</sub>O<sub>5+x</sub> compounds have demonstrated to be very efficient cathode materials not only in solid oxide fuel cells but also in proton conducting fuel cells. In this last case, the excellent properties could be due to the presence of mixed electron-proton conduction. We study here the diffusion of the proton in this material using molecular dynamics simulations. Two different diffusion mechanisms are observed. The predominant mechanism is the standard proton transfer between two neighbouring oxygen atoms combined with the rotation of H around its first neighbour oxygen atom. The second mechanism consists in the migration of the OH group where both oxygen and hydrogen atoms diffuse together. Strong spatial correlations between successive proton jumps are evidenced. This is likely related to the presence of oxygen vacancies and to the concerted diffusion of hydrogen and oxygen atoms.

### 1. Introduction

Since the first article on the use of GdBaCo<sub>2</sub>O<sub>5+x</sub> double perovskite compounds as cathode materials in solid oxide fuel cells [1] and the exhaustive study of their properties by A. Tarancón and co-authors [2-5], the interest for this family of compounds has been growing. This interest principally comes from the excellent properties, in particular the Area Specific Resistance (ASR), observed for this material as an air electrode material. For example, an ASR of 0.25 Ω.cm<sup>2</sup> has been measured at 625°C [2]. Several attempts have been made to improve the chemical stability, the thermal expansion coefficient (too high in the pristine compound) and the electrochemical properties. The main pathways of improvement have consisted in the substitution of Ba by Sr, leading to potentially more stable compounds [6-8], and the use of other transition metals to decrease the thermal expansion [8-11]. A conclusion from these studies is that increasing the chemical stability and lowering the thermal expansion coefficient is possible but most often leads to a decrease of electrochemical performance. Besides these experimental results, theoretical studies have provided a deeper understanding of the atomic scale mechanisms of diffusion in this type of compound. In particular, Molecular Dynamics (MD) simulations have led to the confirmation of the 2D nature of the oxygen diffusion in REBaCo<sub>2</sub>O<sub>5.5</sub> compounds (RE= La, Pr, Gd, Y, ...) and have revealed the associated diffusion mechanisms [12,13]. It has been shown through neutron diffraction analysis combined with MD calculations that long range diffusion occurs through oxygen ion jumps between gadolinium (GdO) and cobalt (CoO<sub>2</sub>) planes, thus emphasising the importance of the presence of oxygen vacancies in the CoO<sub>2</sub> plane [14]. More recently, these double perovskite compounds have been applied in Proton Conducting Fuel Cells (PCFCs), displaying excellent properties [15]. For instance, in La-doped GdBaCo<sub>2</sub>O<sub>5+x</sub>, it was shown using electrode polarization resistance that the apparent activation energy decreases from around 1.3 eV at high temperatures to around 0.5 eV below

550°C [16]. This was attributed essentially to a change from oxygen ion transport at high temperatures to mainly proton transport at low temperatures. This assumption was supported by thermogravimetric measurements showing a weight increase under wet atmospheres for the La-doped compounds. As shown in a recent article, PrBaCo<sub>2</sub>O<sub>5+x</sub> used as electrode material presents a much lower ASR in wet atmosphere compared to dry atmosphere [17]. This result is potentially linked with the incorporation of water inside the double cobaltite compound, incorporation that would favor the oxygen reduction reaction, even if there is no direct proof of proton incorporation in this material. A previous study also showed that REBaCo<sub>2</sub>O<sub>5+x</sub> double perovskites are stable in humid atmospheres and potentially incorporate water at low temperatures [18]. Nevertheless, recent neutron diffraction experiments did not confirm the presence of water in humidified NdBaCo<sub>2</sub>O<sub>5+x</sub> at high temperatures [19]. However, DFT calculations found either exothermic or endothermic hydration of GdBaCo<sub>2</sub>O<sub>5.5</sub> depending on the original structure considered [20]. These calculations also seem to indicate a preferential location of protons close to CoO<sub>2</sub> and GdO planes [21]. Further experimental and theoretical studies are thus needed to investigate the possible hydration of double perovskite cobaltite compounds and to understand its effects on transport and electrochemical properties.

Keeping in mind that the hydration of a double cobaltite compound is still the object of debate, we assume in the present study a partial hydration, and performed MD simulations to understand proton diffusion mechanisms in GdBaCo<sub>2</sub>O<sub>5.5</sub> and the role of hydration on transport properties.

### 2. Computation details

For the present study, the interatomic potentials within GdBaCo<sub>2</sub>O<sub>5.5</sub> are those previously used to study oxygen diffusion in this material [12,13]. These potentials have proven their validity for reproducing oxygen sublattice dynamics [12,13]. Since proton transport is essentially a phonon-assisted process, through oxygen vibrations, all the ingredients are

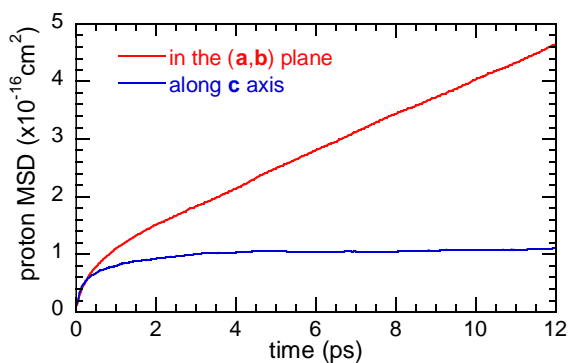
\* Corresponding author: Fabien Briec  
Email address: fabien.briec@centralesupelec.fr

present for reproducing the proton diffusion in double cobaltites. The Buckingham potentials between the first neighbouring oxygen ion of the proton and the other cations were adapted to take into account the smaller charge on the oxygen atom of the hydroxyl group, following a procedure presented by Sierka *et al.* [21]. The proton transfer between the two first neighbouring oxygen atoms is described by the empirical valence bond (EVB) model, as originally presented by Raiteri *et al.* for proton conductors [22,23]. In particular, using this model, the proton diffusion coefficient is correctly reproduced in Y-doped barium zirconate as compared to experimental values [22]. The EVB parameters associated with the interactions between the proton and the other ions were taken from the original article on Y-doped BaZrO<sub>3</sub> [22]. All the parameters of the interatomic potentials are provided as supplementary information.

The size of the MD simulation box was  $6a \times 6b \times 3c$  as compared to the original double perovskite cell, containing 108 Ba, 108 Gd, 216 Co and 594 O, corresponding to the formula GdBaCo<sub>2</sub>O<sub>5.5</sub> (GBCO). In the initial configuration, the GdO plane is deficient in oxygen atoms with respect to the stoichiometric double perovskite structure. The associated empty sites represent oxygen vacancies, which can migrate during the MD simulations. One proton is introduced in the simulation box and the electrical neutrality is ensured by adding a uniform background charge. The simulations were performed with a time step of 0.3 fs in the NVT ensemble using a Langevin thermostat [24] and at a pressure close to zero. The temperature was varied between 1200 and 1900 K and average quantities were obtained using equilibrium trajectories of 11 ns to 2 ns, respectively. All the computations were carried out with a homemade code.

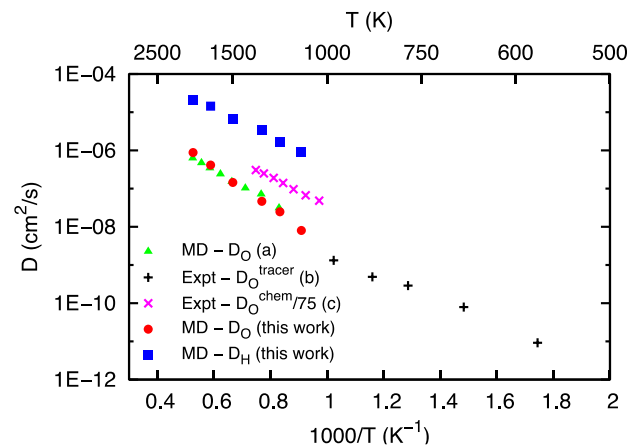
### 3. Results

Simulations first indicate that the presence of one proton out of 1026 atoms does not affect the volume of the cell. As shown in previous studies [12,13,25], oxygen atoms diffuse parallel to the  $(a,b)$  plane but are blocked in the  $c$  direction. Up to 1700 K, BaO planes act as barriers for the diffusion of oxygen atoms, which explains their limited displacement in the  $c$  direction and thus the 2D nature of the diffusion. The oxygen diffusion pathway is composed of a sequence of jumps between the GdO plane and an adjacent CoO<sub>2</sub> plane, similarly to what have been previously obtained in GBCO [12,13,25] and observed experimentally in a similar compound (NdBaCo<sub>2</sub>O<sub>5+x</sub>) [14,26].



**Figure 1.** Mean Square Displacement (MSD) of a proton in GBCO at 1700 K.

The Mean Square Displacement (MSD) of the proton as obtained from simulations at 1500 K is given in figure 1. The proton diffuses in parallel to the  $(a, b)$  plane and is blocked in



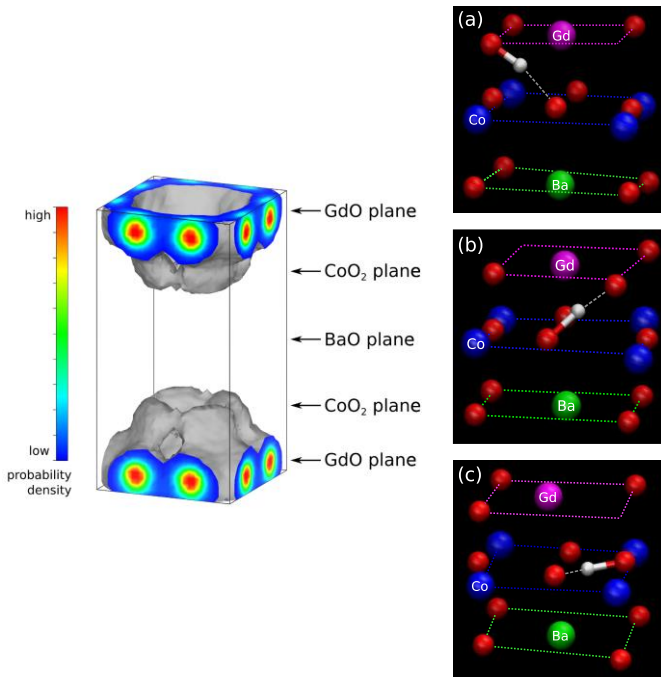
**Figure 2.** Proton (full squares) and oxygen (full circles) diffusion coefficients in the  $(a,b)$  plane,  $D$  (in log scale), as obtained from MD simulations. Previous oxygen results in GBCO are also displayed: (a) MD simulations [12], (b) from oxygen exchange [3] and (c) from conductivity relaxation [27, 13].

the  $c$  direction, similarly to the oxygen atoms. Figure 2 shows the diffusion coefficient in the  $(a, b)$  plane of both protons and oxygen atoms as extracted from MSD. Comparison is made with previous experimental and theoretical results, which are only available for oxygen. The calculated oxygen diffusion coefficient is in reasonable agreement with experiments [3]. For the first time, the proton diffusion coefficient (blue squares in figure 2) is evaluated in GBCO. It is found to be more than one order of magnitude higher than that of oxygen. Furthermore, the activation energy associated with proton diffusion ( $0.71 \pm 0.03$  eV) is found to be lower than the activation energy for oxygen diffusion ( $1.01 \pm 0.03$  eV). This difference of activation energy between oxygen and proton migration is commonly observed in perovskite [28-32] or other oxide compounds [33,34]. For instance, in Y-doped barium zirconate, an activation energy of 0.95 eV [31] is reported for oxygen ion migration, while it is 0.40-0.45 eV [28-32] for the proton migration.

Unfortunately, no data are available concerning proton diffusion/conductivity in double perovskite cobaltites which could be compared to our simulation results. This comes from the fact that proton incorporation through water uptake is difficult to evaluate in electrode materials due to the concomitant oxygen uptake (oxidation reaction). We consider that one method could nevertheless be used, namely nuclear reaction analysis, which would allow the quantification of proton presence inside dense GBCO materials, for concentrations varying from traces to a few percent [35]. This method would not have the drawbacks of thermogravimetry measurements and could give a direct proof of proton incorporation. If the incorporation of protons is confirmed, the proton diffusion coefficient could be evaluated through methods like TOF-SIMS (Time-of-Flight Secondary Ion Mass Spectrometry) which has been applied with success in lanthanum tungstates [36], keeping in mind that this may provide information on water diffusion rather than on proton diffusion.

Before the analysis of the proton migration mechanisms, let us consider its preferential locations in GBCO. The density map on Figure 3 (left) indicates that the proton is essentially located slightly above or below the GdO plane and to a less extent close to the CoO<sub>2</sub> plane. The proton is likely to be rejected from the BaO plane just like oxygen vacancies. Indeed, at 1500 K, 97% of the oxygen vacancies are, at equilibrium, located in GdO planes, while the remaining 3% are in CoO<sub>2</sub> planes. In other words, the proton is never bonded to an oxygen atom of

the BaO plane. The snapshots of Figure 3 (right), represents the preferential locations of the proton. Most of the time (87% at 1500 K), the proton is bonded to an oxygen of the GdO plane (snapshot (a)) while the rest of the time it is bonded to an atom of the CoO<sub>2</sub> plane (snapshots (b) and (c)).



**Figure 3.** Left - Probability density map of the proton in GBCO at 1500K, obtained from MD trajectories. Right – Snapshots of the preferential locations of the proton in GBCO (*ie.* configurations associated with the probability density maximums).

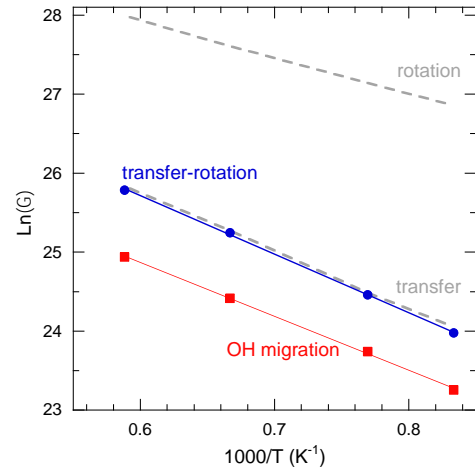
From the density map of figure 3, the proton diffuses within the same GdO and CoO<sub>2</sub> planes since the probability density is zero near the in-between BaO plane. This corroborates the bidimensional nature of proton diffusion parallel to the (*a,b*) plane (see figure 1). The diffusion of protons in GBCO follows a scheme where the proton is successively bonded to one oxygen atom of the GdO plane and then to an oxygen atom of the CoO<sub>2</sub> plane. Moreover, the OH bond most of the time points in a direction that keeps the proton between the GdO and the CoO<sub>2</sub> planes. This explains the high density of protons observed between these two planes.

The analysis of the MD trajectories shows that the long-range migration of the proton is obtained through two mechanisms.

i) The proton transfer between the two neighbouring oxygen atoms (Grotthuss mechanism) combined with the rotation of H around its first neighbour oxygen atom. The limiting step is the proton transfer as shown in figure 4. The activation energies are equal to 0.63 and 0.38 eV for the proton transfer and rotation, respectively. Accounting for the two types of events, the time of flight of the transfer-rotation mechanism is almost fully given by that of the transfer step which is around 0.1 ps. In addition, proton transfer between GdO and CoO<sub>2</sub> planes represents around 93% of the total number of transfer events at T= 1500 K, the others being within the CoO<sub>2</sub> plane (7%).

ii) The second mechanism consists in the migration of the whole OH group (vehicle mechanism [32]) which occurs mainly along <101> directions of the cubic cell. The corresponding time of flight is about 0.3 ps. At 1500K, the contributions of the two diffusion mechanisms, transfer-rotation and OH migration, are 70% and 30%, respectively. Figure 4 shows the Arrhenius plot of their associated jump frequencies. The extracted

activation energies are very similar, being equal to 0.64, and 0.59 eV for the transfer-rotation and OH migration mechanisms, respectively.



**Figure 4.** Arrhenius plot of the frequency,  $\Gamma$  (Hz), of the transfer-rotation (full circles), and OH migration (full squares) mechanisms associated to the diffusion of the proton. Pure rotation and transfer events are also shown (grey dashed lines)

These activation energies are much smaller than the ones calculated from the diffusion coefficient values obtained using the mean square displacement. This can be explained by the presence of spatial correlations between successive jumps leading to a proton diffusion that cannot be described by a purely random walk. The deviation from a purely Brownian diffusion can be quantified using the correlation factor,  $f$  [37,38], which is estimated using the jump frequencies of the transfer-rotation ( $\Gamma_H$ ) and of the OH-migration ( $\Gamma_{OH}$ ) events. In contrast to a purely random walk, for which  $f = 1$ , the jump directions in GBCO may not be equally probable and may depend on the prior jump. The jump sequence is thus correlated and  $f \neq 1$ .

The proton diffusion in GBCO is considered in the (*a, b*) plane since the H atom never visits the BaO plane. The correlation factor is defined as the ratio between the MSD diffusion coefficient,  $D_H$ , and the diffusion coefficient estimated from the jump frequencies in the purely random walk regime:

$$f = \frac{D_H}{\frac{1}{4}(\Gamma_H d_H^2 + \Gamma_{OH} d_{OH}^2)}$$

where the factor 1/4 comes from the 2D nature of the diffusion. The jump frequency  $\Gamma_H$  corresponds to the transfer-rotation combination [39]:

$$\frac{1}{\Gamma_H} = \frac{1}{\Gamma_{rotation}} + \frac{1}{\Gamma_{transfer}}$$

and  $\Gamma_{OH}$  is associated to the OH migration mechanism where both oxygen and hydrogen atoms migrate together. The jump lengths  $d_H$  and  $d_{OH}$  are the proton displacements projected in the (*a, b*) plane. Both lengths are about  $a/2$ , since the jumps are oriented mainly along the [101] or [011] directions in the cubic cell. Our MD values for the correlation factor are equal to 0.28, 0.59, and 0.76 at 1300 K, 1500 K, and 1700 K, respectively. These values less than unity mean that there exists a strong spatial correlation between the jumps. This is likely due to the presence of oxygen vacancies and to the concerted diffusion of oxygen and hydrogen atoms (vehicle diffusion). Indeed, we obtained a correlation factor close to 1

(0.94) in the case of a proton diffusion in BaZrO<sub>3</sub> where the oxygen atoms do not migrate. This temperature-dependent correlation factor may explain in part the high activation energy of the proton migration observed here.

Besides giving some features of the proton diffusion mechanisms, these results provide some evaluation of the potential influence of hydration of GBCO on its transport properties. Indeed, if we consider for instance the filling of 5% of the oxygen vacancies in GdBaCo<sub>2</sub>O<sub>5.5</sub>, we would have a conductivity for protons at 600°C that is 0.5 times that of oxygen ions. In other words, a slight incorporation of protons in GBCO would result in a high proton conductivity at the working temperature of PCFCs *i.e.* ~ 600°C.

#### 4. Conclusions

Using a reactive-force field, we have been able to treat the potential diffusion of protons in GBCO. We have shown that protons would present a 2D diffusion, analogous to that of oxygen in this material. We evaluated the proton diffusion coefficient being more than one order of magnitude greater than the oxygen diffusion coefficient. We also show that the activation energy of the 2D proton diffusion is equal to 0.71 eV, smaller than for the oxygen atoms. By analysing the trajectories, we found that protons jump from oxygen atoms in the GdO plane to oxygen atoms in the CoO<sub>2</sub> plane and reciprocally. The OH bond points in a direction that keeps the proton between the GdO and CoO<sub>2</sub> planes. This corresponds to the prevailing diffusion mechanism where a proton transfer (Grotthuss mechanism) between the two neighbouring oxygen atoms is combined with a rotation of H around its first neighbour oxygen atom. A second mechanism (vehicle mechanism) consists in the migration of the OH group where both oxygen and hydrogen atoms migrate together, mainly along <101> directions of the cubic cell. These two mechanisms have similar activation energies around 0.6 eV. We also evidenced a strong spatial correlation of the atomic jumps, since a correlation factor less than unity was found in the temperature range 1300-1700 K. This correlation seems to be linked to the concerted diffusion of hydrogen and oxygen atoms.

These results are important because they predict a high proton diffusion coefficient in hydrated double perovskite cobaltite compounds, which would strongly help the cathode reaction in PCFCs. Further experiments are nevertheless needed to quantify the hydration degree of these compounds. We estimate that a partial filling by water of 5% of the oxygen vacancies in GBCO would lead to a proton conductivity at 600 °C comparable to that of oxygen ions. For further work, a detailed comparison of the proton diffusion in double perovskite compounds, such as GBCO, and in more standard perovskites, such as BaZrO<sub>3</sub>, could give important information on the differences between these two classes of materials.

#### Acknowledgements

This work was performed using HPC resources from the *mesocentre* computing center of CentraleSupélec, which is supported by CentraleSupélec and CNRS.

#### References

[1] A. Chang, S.J. Skinner, J.A. Kilner, *Solid State Ionics* 177 (2006) 2009-2011  
[2] A. Tarancón, A. Morata, G. Dezanneau, S.J. Skinner, J.A. Kilner, S. Estradé, F. H. Hernandez-Ramírez, F. Peiró, J.R. Morante, *J. Power Sources* 174 (2007) 255-263

[3] A. Tarancón, S.J. Skinner, F. Hernandez-Ramírez, J.A. Kilner, *J. Mater. Chem.* 17 (2007) 3175-3181  
[4] A. Tarancón, D. Marrero-López, J. Peña-Martínez, J.C. Ruiz-Morales, P. Núñez, *Solid State Ionics* 179 (2008) 611-618  
[5] A. Tarancón, J. Peña-Martínez, D. Marrero-López, A. Morata, J.C. Ruiz-Morales, P. Núñez, *Solid State Ionics* 179 (2008) 2372-2378  
[6] H. Ding, X. Xue, *J. Alloys Compd.* 496 (2010) 683-686  
[7] X. Zhang, M. Jin, *J. Power Sources* 195 (2010) 1076-1078  
[8] J. H. Kim, J.T.S. Irvine *Int. J. Hydrogen Energy* 37 (2012) 5920-5929  
[9] Y. Hu, C. Bogicevic, Y. Bouffanais, M. Giot, O. Hernandez, G. Dezanneau, *J. Power Sources* 242 (2013) 50-56  
[10] A. Kulka, Y. Andrzej, G. Dezanneau, J. Molenda, *Funct. Mater. Lett.* 4 (2011) 157-160  
[11] X.L. Che, Y. Shen, H. Li, T. He, *J. Power Sources* 222 (2013) 288-293  
[12] J. Hermet, G. Geneste, G. Dezanneau, *Appl. Phys. Lett.* 97 (2010) 174102  
[13] J. Hermet, B. Dupé, G. Dezanneau, *Solid State Ionics* 216 (2012) 50-53  
[14] Y. Hu, O. Hernandez, T. Broux, M. Bahout, J. Hermet, A. Ottochian, C. Ritter, G. Geneste, G. Dezanneau, *J. Mater. Chem.* 22 (2012) 18744-18747  
[15] B. Lin, S. Zhang, L. Zhang, L. Bi, H. Ding, X. Liu, J. Gao, G. Meng, *J. Power Sources* 177 (2008) 330-333  
[16] R. Strandbakke, V.A. Cherepanov, A.Y. Zuev, D.S. Tsvetkov, C. Argiris, G. Sourkouni, S. Prünke, T. Norby, *Solid State Ionics* 278 (2015) 120-132  
[17] A. Grimaud, J.-M. Bassat, F. Mauvy, M. Pollet, A. Wattiaux, M. Marrony, J.-C. Grenier, *J. Mater. Chem. A* 2 (2014) 3594-3604  
[18] G. Goupil, T. Delahaye, B. Sala, F. Lefebvre Joud, G. Gauthier, *Solid State Ionics* 263 (2014) 15-22  
[19] M. Bahout, S.S. Pramana, J.M. Hanlon, V. Dorcet, R.I. Smith, S. Paofai, S.J. Skinner, *J. Mater. Chem. A* 3 (2015) 15420-15431  
[20] E. Coulaud, G. Dezanneau, G. Geneste, *J. Mater. Chem. A* 3 (2015) 23917-23929  
[21] M. Sierka, J.Sauer, *Faraday Discussions* 106 (1997) 41-62  
[22] P. Raiteri, J.D. Gale, G. Bussi, *J. Phys.: Condens. Matter* 23 (2011) 334213  
[23] A. Ottochian, G. Dezanneau, C. Gilles, C. Knight, P. Raiteri, J.D. Gale, *J. Mater. Chem. A* 2 (2014) 3127-3133  
[24] M. Binggeli, J. R. Chelikowsky, *Phys. Rev. B* 50 (1994) 11764  
[25] D. Parfitt, A. Chronos, A. Tarancón, J. A. Kilner, *J. Mater. Chem.* 21 (2011) 2183-2186  
[26] R. A. Cox-Galhotra, A. Huq, J. P. Hodges, J.-H. Kim, C. Yu, X. Wang, A. J. Jacobson and S. McIntosh, *J. Mater. Chem. A* 1 (2013) 3081-3100  
[27] M.-B. Choi, S.-Y. Jeon, J.-S. Lee, H.-J. Hwang, S.-J. Song, *J. Power Sources*, 195 (2010) 1059-1064  
[28] K. Wakamura, *J. Phys. Chem. Solids* 66 (2005) 133-142  
[29] S. Ricote, N. Bonanos, A. Manerbino, W.G. Coors, *Int. J. Hydrogen Energy* 37 (2012) 7954-7961  
[30] S. Ricote, N. Bonanos, H.J. Wang, B.A. Boukamp, *Solid State Ionics* 213 (2012) 36-41  
[31] A. C. T. van Duin, B. V. Merinov, S. S. Han, C. O. Dorso, W. A. Goddard III, *J. Phys. Chem. A* 112 (2008) 11414-11422  
[32] K.-D. Kreuer, *Annu. Rev. Mater. Res.* 33 (2003) 333-359  
[33] Y. Larring, T. Norby, *Solid State Ionics* 77 (1995) 147-151  
[34] Y. Larring, T. Norby, *Solid State Ionics* 97 (1997) 523-528  
[35] S. Sorieul, S. Miro, M. Taillades-Jacquín, J. Dailly, F. Mauvy, M.-H. Berger, P. Berger, *Nucl. Instr. Meth. Phys. Res. B* 266 (2008) 1430-1433

*Solid State Ionics*, Elsevier, 2017, 309, pp.187-191. [10.1016/j.ssi.2017.07.017](https://doi.org/10.1016/j.ssi.2017.07.017)

[36] R. Hancke, S. Fearn, J.A. Kilner, R. Haugrud, Phys. Chem. Chem. Phys. 14 (2012) 13971-13978  
[37] H. Mehrer, Diffusion in solids, Springer series in solid state science, Springer, Berlin and Heidelberg, 2007

[38] J. Philibert, Atom movements diffusion and mass transport in solids, Les Editions de physique, Les Ulis, 1991  
[39] M. E. Björketun, P. G. Sundell, G. Wahnström, D. Engberg, Solid State Ionics 176 (2005) 3035-3040

## Supplementary Information

### 1. Interatomic potential

The charges for cations are +3, +2 and +3 for Gd, Ba and Co, respectively. O refers to oxygen ions, while O<sub>OH</sub> refers to an oxygen ion bonded to a proton. The charge of O, O<sub>OH</sub> and OH group are -2, -1.308698 and -1, respectively. The values of the parameters of the Buckingham potentials used for molecular dynamics calculations in GBCO are given in table 1.

Interaction	A (eV)	$\rho$ (Å)	C (eV.Å <sup>6</sup> )	Ref.
Gd-O	1458.38	0.3522	0.0	[a]
Ba-O	1214.4	0.3537	0.0	[a,b]
Co-O	1329.82	0.3087	0.0	[a,b]
O-O	22764.3	0.149	43.0	[a,b]
Gd-O <sub>OH</sub>	856.71	0.3522	0.0	this work
Ba-O <sub>OH</sub>	1214.3	0.3537	0.0	this work
Co-O <sub>OH</sub>	1137.66	0.3087	0.0	this work
O-O <sub>OH</sub>	22764.3	0.149	43.0	this work
O <sub>OH</sub> -O <sub>OH</sub>	22764.3	0.149	43.0	this work

**Table 1.** Parameters of the Buckingham potential.

a : J. Hermet, G. Geneste, G. Dezanneau, *Applied Physics Letters*, 2010, **97**, Article N° 174102

b : C. A. J. Fisher, M. Yoshiya, Y. Iwamoto, J. Ishii, M. Asanuma, and K. Yabuta, *Solid State Ionics*, 2007, **177**, 3425

### 2. The Empirical Valence Bond (EVB) model

Proton reactivity is included through the EVB model proposed by Raiteri and coworkers [P. Raiteri, J.D Gale, G. Bussi, *J. Phys.: Condens. Matter*, 2011, **23**, 334213]. In this model, the proton is considered to be in two different states where it is bonded to different oxygen atoms. State 1: H is bonded to the nearest oxygen atom O<sub>1</sub> with associated potential energy V<sub>1</sub>; State 2: H is bonded to the second nearest oxygen atom O<sub>2</sub> with associated energy V<sub>2</sub>. These two states are coupled by the term V<sub>12</sub>, so that the total potential energy can be written in a matrix form:

$$V = \begin{pmatrix} V_1 & V_{12} \\ V_{12} & V_2 \end{pmatrix}$$

V<sub>12</sub> is given by:

$$V_{12} = \lambda \exp(-\xi Q^2)$$

where Q is the reaction coordinate associated to the transfer of the proton between O<sub>1</sub> and O<sub>2</sub>

$$Q = d_{O_2H} - d_{O_1H}$$

$d_{O_2H}$ ,  $d_{O_1H}$  being the distances between the proton and the O<sub>2</sub> and O<sub>1</sub> atoms, respectively.

The parameters  $\lambda = 0.7998$  eV and  $\xi = 16.0$  Å<sup>-2</sup> comes from ref. [22].

Finally, the potential energy of the system including a hydrogen atom is given by the lowest eigenvalue of the V matrix:

$$V = \frac{(V_1 + V_2) - \sqrt{(V_1 - V_2)^2 + 4V_{12}^2}}{2}$$