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Magnetoelastic standing waves induced in UO_2 by microsecond magnetic field pulses

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Magnetoelastic measurements in the piezomagnetic antiferromagnet UO_2 were performed via the fiber Bragg grating method in magnetic fields up to 150 T generated by a single-turn coil setup. We show that in short timescales, order of a few micro seconds, pulsed-magnetic fields excite mechanical resonances at temperatures ranging from 10 K to 300 K, in the paramagnetic as well as within the robust antiferromagnetic state of the material. These resonances, which are barely attenuated within the 100 ms observations, are attributed to the strong magnetoelastic coupling in UO_2 combined with the high crystallographic quality of the single crystal samples. They compare well with mechanical resonances obtained by a resonant ultrasound technique and superimpose on the known non-monotonic magnetostriction background. A clear phase-shift of π in the lattice oscillations is, unexpectedly, observed in the antiferromagnetic state when the magnetic field overcomes the piezomagnetic switch-field $H_c \simeq -18$ T. We further present simulations and a theoretical argument to explain the observed phenomena.

I. INTRODUCTION

The antiferromagnetic (AFM) insulator Uranium dioxide UO_2 has been the subject of extensive research during the last decades predominantly due to its widespread use as nuclear fuel in pressurized heavy water reactors. Besides efforts to understand the unusually poor thermal conductivity of UO_2 which impacts its performance as nuclear fuel [1], a recent magnetostriction study in pulsed magnetic fields to 92 T uncovered linear magnetostriction in UO_2 [2] - a hallmark of piezomagnetism.

Piezomagnetism is characterized by the induction of a magnetic polarization by application of mechanical strain, which, in the case of UO_2 , is enabled by broken time-reversal symmetry in the 3- k antiferromagnetic structure which emerges below $T_N = 30.8$ K [3–6] and is accompanied by a Jahn-Teller distortion of the oxygen cage [7–10]. This also leads to a complex hysteretic magnetoelastic memory behavior where magnetic domain switching occurs at fields around ± 18 T at $T = 2.5$ K. Interestingly, the very large applied magnetic fields proved unable to suppress the AFM state that sets in at T_N [2]. These earlier results provide direct evidence for the unusually high energy scale of spin-lattice interactions, and call for further studies in higher magnetic fields.

Here we present axial magnetostriction data obtained in a UO_2 single crystal in magnetic fields to 150 T.

These ultra-high fields were produced by single-turn coil pulsed resistive magnets and applied along the [111] crystallographic axis at various temperatures between 10 K and room temperature. We see, at all temperatures, a dominant negative magnetostriction proportional to the square of the applied field accompanied by unexpectedly strong oscillations that establish a mechanical resonance in the sample virtually instantly upon delivery of the ultra-fast, 10^2 T/ μs , magnetic field rate-of-change. The oscillations observed quickly set, well within a single oscillation period, are long-lasting due to very low losses, with frequencies in the hundreds of kilohertz that match mechanical resonances obtained with a resonant ultrasound spectroscopy (RUS) technique [11]. When the sample temperature is reduced, the frequencies soften, consistent with observations in studies of the UO_2 elastic constant c_{44} as a function of temperature [12]. When the magnetic field is applied at temperatures $T < 30.8$ K in the AFM state, the magnetic field changes sign to a negative field magnitude (a characteristic of destructive magnets) in excess of the UO_2 AFM domain-switch-field of $\simeq -18$ T. This negative field that follows a positive field pulse exposes yet another unexpected result, namely a π (180°) phase-shift in the magnetoelastic oscillations. We use a driven harmonic oscillator and an analytical model to shed light on the origin of our findings.

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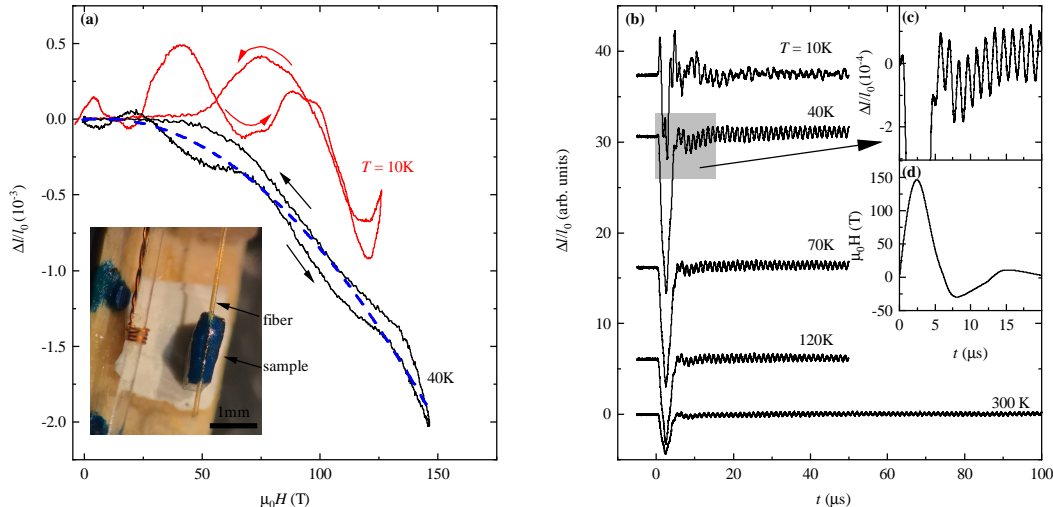


FIG. 1. (a) Uniaxial magnetostriction $\Delta l/l_0$ of a UO_2 single crystal along the $[111]$ axis for $T = 10\text{ K}$ (red line) and 40 K (black line). The dashed blue line is a 2nd order polynomial fit of the 40 K data representing the expected magnetostriction behavior. The arrows marking up and down sweep of the magnetic field. The inset shows a picture of the UO_2 single crystal attached to the FBG-furnished fiber. Note that the sample is not physically attached to the sample holder, other than through the $125\ \mu\text{m}$ optical fiber. A copper coil, located next to the sample was used to measure the magnetic field. (b) $\Delta l/l_0$ of UO_2 as a function of time for pulses at different temperatures between 10 and 300 K . The spike in $\Delta l/l_0$ during the first few microseconds is caused by the magnetostriction of the sample (shown in (a)) due to the changing magnetic field during the field pulse (d), whereas the oscillatory part of the signal continues after the field decayed to zero. A detailed view of the oscillations at 40 K is shown in the inset (c) as well as the magnetic field vs. time during the first $20\ \mu\text{s}$ (d).

II. METHODS

In this work, the magnetostriction signal of UO_2 was measured with a 100 MHz coherent pulse fiber Bragg interrogation method. The setup is driven by a mode-locked pulsed 90 fs Er laser with a 100 MHz repetition rate and allows interrogation speed on the 10 ns scale. This method offers a faster readout rate than traditional fiber Bragg grating (FBG) interrogation systems which operate in the range of several kHz [13]. The UO_2 single crystal was attached to the optical fiber using an epoxy encapsulant with the crystallographic $[111]$ axis aligned parallel to the fiber and the magnetic field. A picture of the sample is shown as an inset in figure 1(a). Details about the FBG setup can be found in Ref. [14]. During the field pulse, the induced voltage in a small copper coil, located in close proximity to the sample (see picture in figure 1(a)), was used to measure the magnetic field.

The magnetic field was generated with a semi destructive capacitor-driven single-turn coil magnet system at the National High Magnetic Field Laboratory's pulse field facility at Los Alamos National Laboratory. The system is designed for fields up to 300 T with a rise time of approximately $2.5\ \mu\text{s}$. Note that for magnetic fields in the region of 200 T and above, damage to the cryostat and sample becomes increasingly likely. Therefore magnetic fields were limited to $\approx 150\text{ T}$ in this study,

with a peak rate of change in the order of $10^2\ \text{T}/\mu\text{s}$. Further details about the single-turn coil setup are presented in Refs. [15, 16]. Optical measurement techniques, like the FBG method used here, are in general advantageous in single-turn experiments when compared to, *e.g.*, electrical capacitance-based dilatometry measurements predominantly due to optical fibers being impervious to the large induced voltages generated inside even small metallic loops caused by the large dB/dt , as well as the associated electromagnetic noise.

The measurement of the natural mechanical resonances of elastic vibration where several normal modes of the sample are determined, is obtained with a set of piezoelectric transducers using a technique known as resonant ultrasound spectroscopy. Here, one transducer serves as source of the tunable sinusoidal wave of frequency f and the other serves as detector at the synchronous frequency of the sample's response. The electronics and room temperature apparatus was described in detail by Balakirev *et al.* [11]. In our case, the transducer had an Al_2O_3 hemisphere that allows precise and reproducible point contact on desired positions of the crystal [17]. As a frequency scan is performed, a resonance peak is observed at each of the normal modes. We performed resonant mode measurements on UO_2 single crystals alone as well with the $125\ \mu\text{m}$ optical fiber attached.

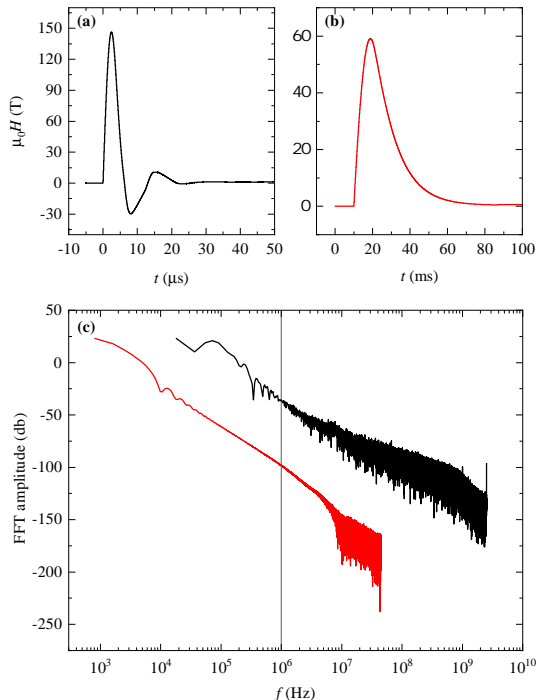


FIG. 2. Magnetic field vs. time of a 146 T pulse in a single-turn coil (a) and a 60 T pulse in a non-destructive short pulse magnet (b). (c) Fourier transformations of the field vs. time curves shown in (a, b). FFT's of the short pulse magnet/single-turn pulses are depicted by red/black lines respectively.

III. RESULTS

Magnetostriction $\Delta l/l_0$ vs. magnetic field curves at $T = 10$ K and 40 K are displayed in Fig. 1. We observe an overall negative magnetostriction signal at high fields as expected from previous results in pulsed magnetic fields to 92 T [2] with no indication of suppression of the robust AFM order. However, the signal is highly hysteretic due to the large mechanical resonances that are superimposed on the magnetostriction signal. The magnetostriction signal itself roughly follows a second order polynomial field dependence (blue dashed line).

The field-induced mechanical resonances become clearer when $\Delta l/l_0$ is plotted as a function of time (Fig. 1(b)). Oscillations start with the onset of the field pulse and persist during the entire data acquisition period (100 μ s at 300 K, 50 μ s for lower temperatures). The high quality of our UO_2 crystal is probably a key factor behind the low attenuation of the mechanical resonances observed in the experiment. This is validated by a large quality factor $Q = f/w$ found by RUS and ranging from 2500 to 5000. Here, f is the frequency and w the width of the resonance (fitted by a Lorentzian). The onset of the mechanical resonances is approximately instantaneous, which indicates that they arise as a response to

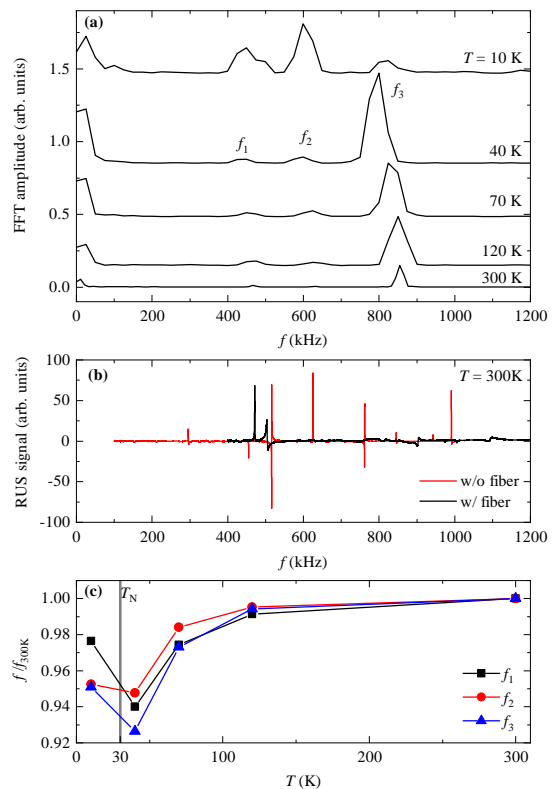


FIG. 3. (a) Fourier transforms of $\Delta l/l_0$ shown in figure 1(b) displaying three dominant frequencies labeled as f_1 , f_2 and f_3 . The low frequency peaks originate from the background magnetostriction. (b) Resonant ultrasound spectroscopy (RUS) spectra at $T = 300$ K. The RUS spectra of the bare sample (black line) and the sample with the attached optical fiber (red line) were recorded along the [111] axis. (c) Normalized temperature dependence of $f_{1,2,3}$.

the magnetic field change and the strong magnetoelastic coupling. Hence, it does not appear that this mechanism is triggered by the shock wave generated by the disintegration of the single turn coil which would need a few microseconds to reach the sample. A similar experiment run with identical interrogation parameters and a bare FBG sensor, i.e. with no sample attached to the fiber, yielded no detectable mechanical resonances.

In order to understand the origin of the observed mechanical resonances, we compare the typical field vs. time profile of a 146 T pulse performed in a single turn coil (Fig. 2(a)) with a 60 T shot in a short-pulse magnet (Fig. 2(b)). The field generated by the short pulse magnet has a total duration of about 100 ms with a rise time of 10 ms. The single turn coil on the other hand has a pulse duration in the order of 25 μ s and a rise time of 2.5 μ s ($\times 4000$ faster compared to the non-destructive magnet). Furthermore, the field switches sign several times, referred to as magnetic field recoil, displaying a significantly less attenuated behavior than the short-pulse magnet. The extremely short timescales in the single turn pulse re-

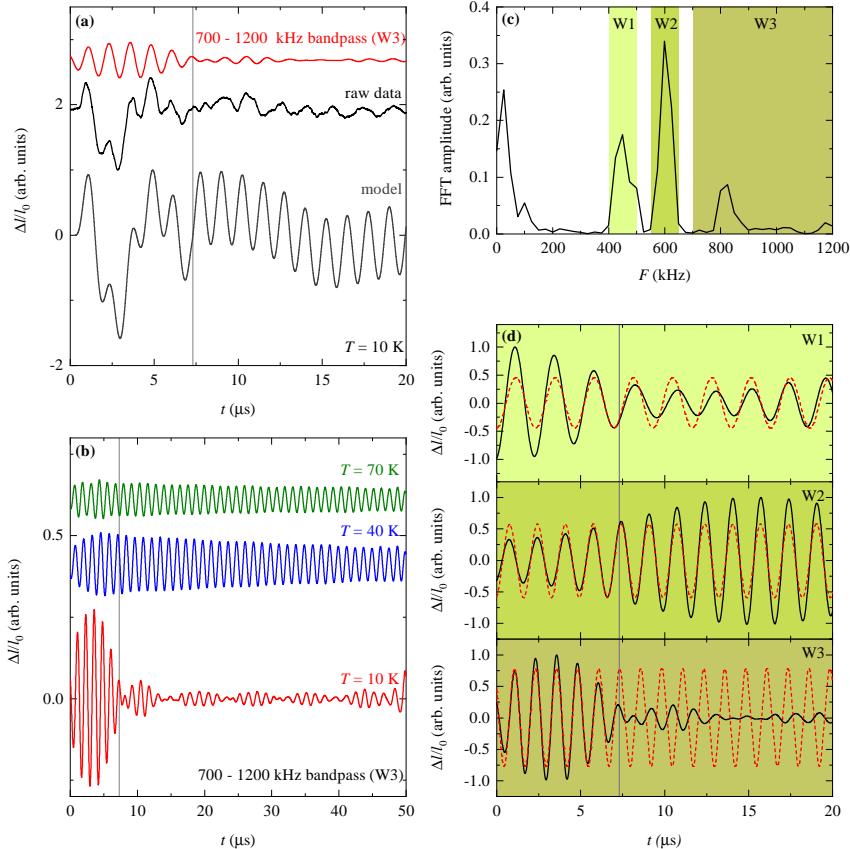


FIG. 4. (a) $\Delta l/l_0$ vs. time at 10 K with a peak field of 126 T before (black line) and after (red line) the 700-1200 kHz band pass filter was applied. The experimental data is compared to a driven harmonic oscillator model with a characteristic frequency of 800 kHz matching the dominant oscillations seen during the field pulse in the magnetostriction signal (grey curve). The experimental data and the model show a π phase shift around $7.3 \mu\text{s}$ marked by the grey vertical line. (b) Magnetoelastic oscillations at 10, 40 and 70 K after application of a 700-1200 kHz band pass filter to the experimental data. Curves are shifted for clarity in figures (a, b). (c) Fast Fourier Transform of $\Delta l/l_0$ at $T = 10$ K. The windows of the band pass filter (W1, W2, W3) are highlighted in different colors. The filtered curves (black lines) for each window are shown in (d) and are compared with sinusoidal functions indicated as red dashed lines.

sults in a shift of the Fast Fourier Transform (FFT) of the field pulse towards higher frequencies up to several MHz (i.e., high frequencies are 50 dB more intense) which is displayed in Fig. 2(c). Thus, if the system under study is magnetic - the field-pulse itself can, by virtue of the strong magnetoelastic coupling present in UO_2 [12, 18, 19], excite mechanical resonances in the range of several 100 kHz.

The FFT of $\Delta l/l_0(t)$ in UO_2 reveals three distinct frequencies labeled f_1 , f_2 and f_3 for temperatures between 40 and 300 K, as shown in Fig. 3(a). All modes display a softening as the temperature is lowered and a stiffening below T_N (Fig. 3(c)), in agreement with previous measurements of the elastic constants of UO_2 which show a similar behavior [12]. Three independent elastic constants c_{11} , c_{12} and c_{44} exist for a lattice with cubic symmetry. The distinct temperature dependence of all three mechanical resonance frequencies f_{1-3} indicates that we

probe predominantly resonances that are associated with c_{44} [12]. The observed frequencies are in agreement with RUS data obtained at 300 K, shown in Fig. 3(b). The RUS measurements reveal a rich spectrum of sharp resonances in the frequency range up to 1 MHz. When the crystal is attached to the fiber, a broadening of resonance peaks is observed, as well as a decrease in amplitude, indicating a larger damping, product of a larger system consisting of the crystal, glue and fiber. Nevertheless, the overall range of mechanical resonance frequencies observed in the magnetostriction measurements is comparable with RUS spectra. In the FBG measurements we observe three distinct frequencies. The additional mechanical resonance frequencies present in the RUS spectra are potential resonances that are either completely damped by the attached fiber and encapsulant or they only have a small compressive component parallel to the fiber since the FBG method is less sensitive to shear strain. Note

that the RUS measurements on the bare UO_2 crystal were performed on a ≈ 2 mm longer sample than the FBG and RUS measurements with the optical fiber attached. The change of the sample geometry also affects the mechanical resonance frequencies.

Upon cooling the sample below the antiferromagnetic transition we observe a substantial change in the mechanical resonances when compared to temperatures above T_N :

(i) In the data recorded at 10 K, within the AFM ordered phase, the resonances appear to be significantly damped after $t \approx 7 \mu\text{s}$, particularly the FFT amplitude of $f_3 = 800$ kHz is suppressed when compared to f_1 and f_2 , in contrast to the Fourier transforms of the data sets above T_N where f_3 is clearly the dominant resonance (Fig. 3(a)). The attenuation effect becomes more apparent when band pass filters are applied to the experimental data, isolating the individual resonances and removing the magnetostriction background (Fig. 4(a, b, d)). We show that, when compared to f_3 , the amplitude of lower resonances f_2 and f_3 does not display a drastic change for $t > 7 \mu\text{s}$. The beating pattern observed in f_3 indicates the presence of two resonances very close to each other, which we cannot easily resolve in the FFT. These details are also impacted by the lower and upper limits chosen for the band pass filter.

(ii) A π phase shift can be observed in the mechanical resonances at $t \approx 7 \mu\text{s}$ which is accompanied by the observed attenuation effect. One can clearly identify the π phase shift after a 700-1200 kHz high pass filter was applied to the experimental data (Fig. 4(d)). Interestingly only f_3 shows the phase shift and both lower resonances f_1 and f_2 seem not to be affected and follow a single sinusoidal function as shown in Fig. 4(d)).

The phase shift around $t \approx 7 \mu\text{s}$ coincides with a magnetic field value of approximately -18 T close to the field value where a abrupt sign change of the AFM ordering vector L_0 (as defined in Ref. [20]) leads to a jump in the lattice distortion and the characteristic piezomagnetic butterfly that was reported in Ref. [2]. This effect is illustrated in Fig. 5 with the piezomagnetic butterfly shown in the inset. The origin of the phase shift can be found in the sudden reversal of L_0 (red) to the $-L_0$ (blue) as we demonstrate in the following section.

IV. HAMILTONIAN

Following Bar'yakhtar *et al.* [20], we denote as L_1 , L_2 and L_3 the different AFM vectors that describe the $3-k$ order in UO_2 , and M the average magnetization. The magnetic unit cell below the AFM transition is made of 4 formula units, each formula unit carrying a magnetic moment S_1 , S_2 , S_3 , and S_4 . The above AFM vectors and magnetization can be expressed, in terms of the individual magnetic moments as:

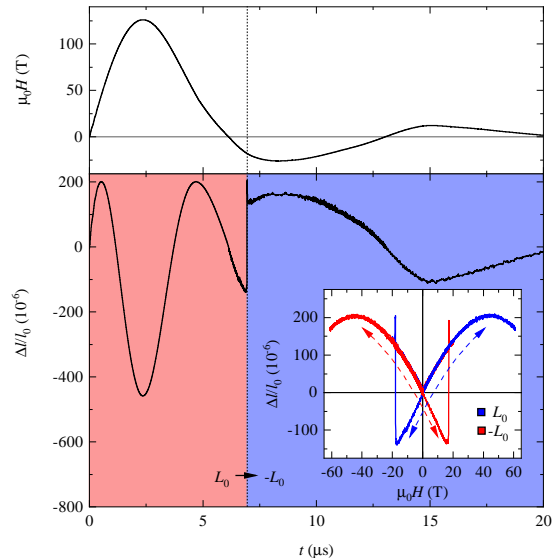


FIG. 5. Magnetic field (upper panel) and magnetostriction of UO_2 (lower panel) as a function of time. The magnetostriction data was mapped and extrapolated onto the single turn field profile using previously published pulse field data (Ref. [2]) shown in the inset. The vertical line at $t \approx 7 \mu\text{s}$ separates the two states with positive (red) and negative (blue) AFM ordering vector L_0 by reaching the switching field of ≈ -18 T in this case.

$$\begin{aligned} M &= S_1 + S_2 + S_3 + S_4, \\ L_1 &= S_1 + S_2 - S_3 - S_4, \\ L_2 &= S_1 - S_2 + S_3 - S_4, \\ L_3 &= S_1 - S_2 - S_3 + S_4, \end{aligned}$$

In the Landau approach, the thermodynamic potential must be able to describe both the paramagnetic and AFM phases. This is achieved by performing a polynomial expansion of the free energy whose terms respect the symmetry of the highest symmetry phase [21]. Such terms have already been worked out in Ref. [20]. The Hamiltonian of the system can thus be written as:

$$\hat{H} = \int d^3\mathbf{r} \hat{\mathcal{H}}(\mathbf{r}), \quad (1)$$

where $\hat{\mathcal{H}}$ can be found in the supplemental material. Since we are primarily interested in explaining the properties of the mechanical resonances and the π phase shift in response to an applied pulsed magnetic fields, we make the simplifying following assumption: the magnetic response of the sample is primarily determined by the external magnetic field, and elastic vibrations do not affect it significantly. As a result, we write each magnetic quantity under the form $\boldsymbol{\mu} = \boldsymbol{\mu}_{eq} + \delta\boldsymbol{\mu}$ (with $\boldsymbol{\mu} = L_1, L_2, L_3, M$). $\boldsymbol{\mu}_{eq}$ represents an equilibrium value,

and $\delta\boldsymbol{\mu}$ the deviation away from that value, *i.e.* the response to the magnetic pulse.

We then write the mechanical equations of motion, from which we retain $\rho\ddot{u}_x$

$$\rho\ddot{u}_x = \frac{\partial}{\partial x} \left(\frac{\partial \hat{\mathcal{H}}}{\partial \eta_1} \right) + \frac{\partial}{\partial y} \left(\frac{\partial \hat{\mathcal{H}}}{\partial \eta_6} \right) + \frac{\partial}{\partial z} \left(\frac{\partial \hat{\mathcal{H}}}{\partial \eta_5} \right), \quad (2)$$

In the mechanical equations, ρ is the volumic mass and u_α is the displacement field in direction α . Since the

derivations are lengthy, we will focus on the component u_x of the displacement field which, after replacing with the Hamiltonian expression (see SI) and assuming, as demonstrated in Ref. [20], that at equilibrium, $L_{1x}^{eq} = L_0$, $L_{2y}^{eq} = L_0$ and $L_{3z}^{eq} = L_0$, we write $L_{1x} = L_0 + \delta L_{1x}$, $L_{1y} = \delta L_{1y}$, $L_{1z} = \delta L_{1z}$, etc. Similarly, Ref. [2] shows that no magnetization seems to exist in the antiferromagnetic phase, so we write $M_x = \delta M_x$, etc. We also recall that $\eta_1 = \frac{\partial u_x}{\partial x}$, $\eta_5 = \frac{\partial u_x}{\partial z} + \frac{\partial u_z}{\partial x}$, the equation can now be written, for linear terms, as

$$\begin{aligned} \rho\ddot{u}_x \approx & c_{11} \frac{\partial^2 u_x}{\partial x^2} + c_{12} \left(\frac{\partial^2 u_y}{\partial x \partial y} + \frac{\partial^2 u_z}{\partial x \partial z} \right) + c_{44} \left(\frac{\partial^2 u_x}{\partial y^2} + \frac{\partial^2 u_y}{\partial y \partial x} + \frac{\partial^2 u_z}{\partial z \partial x} + \frac{\partial^2 u_x}{\partial z^2} \right) \\ & + L_0 (\lambda_1 + \lambda'_1) \frac{\partial \delta L_{1x}}{\partial x} + \lambda'_1 L_0 \left(\frac{\partial \delta L_{2y}}{\partial x} + \frac{\partial \delta L_{3z}}{\partial x} \right) + \frac{\lambda}{2} L_0 \left(\frac{\partial \delta L_{1y}}{\partial y} + \frac{\partial \delta L_{2x}}{\partial y} + \frac{\partial \delta L_{1z}}{\partial z} + \frac{\partial \delta L_{3x}}{\partial z} \right). \end{aligned} \quad (3)$$

In Equation 3, the terms proportional to λ_1 , λ'_1 and λ relate the linear change of shape of an antiferromagnetic uniform domain with respect to antiferromagnetic excitation.

If we perform a Fourier transform $u_x = \int \frac{d\mathbf{q}}{2\pi} \int \frac{d\omega}{2\pi} u_x(\mathbf{q}, \omega) e^{i(\mathbf{q}\cdot\mathbf{r} - \omega t)}$, etc., yielding:

$$\begin{aligned} (c_{11}q_x^2 + c_{44}[q_y^2 + q_z^2] - \rho\omega^2) u_x(\mathbf{q}, \omega) = & -(c_{44} + c_{12})q_x [q_y u_y(\mathbf{q}, \omega) + q_z u_z(\mathbf{q}, \omega)] \\ & + iq_x L_0 [(\lambda_1 + \lambda'_1) \delta L_{1x}(\mathbf{q}, \omega) + \lambda'_1 (\delta L_{2y}(\mathbf{q}, \omega) + \delta L_{3z}(\mathbf{q}, \omega))] \\ & + \frac{\lambda}{2} L_0 (q_y [\delta L_{1y}(\mathbf{q}, \omega) + \delta L_{2x}(\mathbf{q}, \omega)] + q_z [\delta L_{1z}(\mathbf{q}, \omega) + \delta L_{3x}(\mathbf{q}, \omega)]) \end{aligned} \quad (4)$$

We now have sets of coupled harmonic oscillators which are driven by a force which is proportional to L_0 (a cyclic permutation $x \rightarrow y \rightarrow z$ allows to get the equations for the other components). In other words, given a proper change of basis, we can diagonalize this set of equations and write the displacement fields dynamical equations under the form

$$(\omega_0(\mathbf{q})^2 - \omega^2) u_1(\mathbf{q}, \omega) = F_1(\mathbf{q}, \omega),$$

with $F_1(\mathbf{q}, \omega)$ being the force driving the oscillation of u_1 at pulsation ω with wavevector \mathbf{q} . This has an obvious solution, which is

$$u_1(\mathbf{q}, \omega) = \frac{F_1(\mathbf{q}, \omega)}{\omega_0(\mathbf{q})^2 - \omega^2}.$$

We note that F_1 is proportional to L_0 , the antiferromagnetic order parameter. It is then clear that upon reversal of the AFM order, $L_0 \rightarrow -L_0$, the force applied on the set of harmonic oscillators reverses sign, *i.e.* $F_1 \rightarrow -F_1 = e^{i\pi} F_1$ and thus a π phase shift must be experienced in the elastic oscillatory response of the sample long enough after the pulse. Hence, sufficient switching of the AFM order is likely the cause of the phase shift observed in Fig. 4. It is to be noted that the main energy couplings responsible for such an effect are quadratic in the AFM vectors and linear in strain; in other words, they are typical (antiferro)magnetostriction terms. We note that some of those energy couplings are the same ones from which piezomagnetism arises (see Equation 31-32 from Ref. [20]).

Therefore, by assuming that the force driving the oscillations is proportional to the systems strain (shown in Fig. 5) and fixing the frequency at 800 kHz, we are able to model the experimental data with a simple driven harmonic oscillator. As depicted in Fig. 4(a) this harmonic oscillator model reproduces the magnetostriction background as well as the π phase shift in the 800 kHz oscillations. The attenuation observed in the oscillations after the switching of L_0 is not captured by the model and will be discussed in detail below.

V. DISCUSSION

A recent X-ray study on UO_2 single crystals evidences the presence of AFM domains and subsequently the coexistence of AFM phases L_0 and $-L_0$ connected by time-reversal even in magnetic fields beyond the piezomagnetic switching field [22]. This is also supported by magnetostriction measurements [2], which show that the first pulse taken below T_N always has a smaller magnetostriction slope for fields below the switching field. This could be caused by the coexistence of all possible domains, with some contracting and some expanding as the field increases. In our measurements we observe a large attenuation effect around the switching field but oscillations seem not to be further damped afterwards. Therefore, the attenuation of f_3 appears to be caused by the coupling of the mechanical resonances to critical spin fluctuations and/or domain movement close to the piezomagnetic switching field which can lead to a significant attenuation of the mechanical resonances similar to the dramatically increased ultrasonic attenuation that was observed in UO_2 in the vicinity of the AFM phase transition [12]. The mechanical resonances f_1 and f_2 might have a predominantly transversal character which would explain the smaller amplitude and the absence of attenuation below T_N since the longitudinal or compressive modes are expected to be more affected by spin fluctuations [23]. For future experiments we plan to perform magnetocaloric measurements to detect possible heating effects at the switching field caused by dissipative processes like domain movement.

Another interesting point is that the π phase shift only occurs in f_3 . The effect is completely absent in f_2 and much less clear in f_1 which is only slightly out of phase when compared to the single sinusoidal function in the time interval between 0 and 20 μs (Fig. 4(d)). As of now we do not have a conclusive argument on why the phase shift is only visible in f_3 . Depending on the involved antiferromagnetic excitations and the anisotropy of the magneto-elastic couplings, longitudinal and transversal mechanical resonances can display different phase shifts. A possible way to test this in future experiments is to use a birefringent FBG which can yield an orthogonal biaxial strain response along two directions with polarization based probing techniques.

We demonstrate that mechanical resonances can be a

useful tool to detect otherwise-elusive AFM domain flips, and possibly also other types of crystallographic domain dynamics (e.g., in liquid crystals). On the other hand, our results indicate that mechanical resonances can also cause issues in experiments where they are unwanted. A mitigation strategy in experiments where excessive noise is prevalent could consist of clamping the sample as well as to conduct runs with different sample and/or sample-holders geometries and dimensions to minimize mechanical resonances triggered by the magnetic field. This phenomena is reminiscent of wire-motion resonances in electrical transport experiments performed in short-pulse magnets, which can be quite detrimental to the data quality and which effects are minimized by fixing the wires and in this way effectively shifting their resonances to frequencies outside of the experimental range of interest.

VI. SUMMARY

We measured the lattice dilation along [111] for the first time up to 150T in UO_2 , in the AFM as well as in the paramagnetic states. Surprisingly, the AFM state is robust against a field $\geq 120\text{T}$ at 10 K, energy-wise $\sim 4\times$ stronger than $T_N = 30\text{K}$ (if $g = 2$). This result confirms the large energy scale for correlations in UO_2 . We show that mechanical resonances can be induced virtually instantaneously via the magnetoelastic coupling in UO_2 by μs field pulses generated with the single turn coil technique, making this material an interesting candidate for magneto-elastic transducers. We demonstrate the impact of the piezomagnetic switching in UO_2 on the standing wave indicated by a π phase shift and a distinct mode dependent attenuation of the mechanical resonances. Our findings present a novel way to study magnetic dynamics in high magnetic fields and could have an impact on the interpretation of past and future data collected in experiments involving semi destructive pulsed magnetic fields as well practical implications, e.g., as a way to trigger resonators at faster speeds.

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Supplemental Material

I. HAMILTONIAN

As mentioned in the main text, the terms of the Hamiltonian can be derived from the polynomial expansion of the free energy and were worked out by Bar'yakhtar *et al.*. They can be written as:

$$\begin{aligned}
\hat{\mathcal{H}} = & \frac{1}{2}\alpha_L^0(T - T_N) [L_{1x}^2 + L_{2y}^2 + L_{3z}^2] + \frac{1}{4}\beta_L [L_1^4 + L_2^4 + L_3^4] \\
& + \frac{1}{2}\alpha_M M^2 + \frac{1}{2}\lambda_{ML} [(\mathbf{M} \cdot \mathbf{L}_1)^2 + (\mathbf{M} \cdot \mathbf{L}_2)^2 + (\mathbf{M} \cdot \mathbf{L}_3)^2] \\
& + \frac{c_{11}}{2} [\eta_1^2 + \eta_2^2 + \eta_3^2] + c_{12} [\eta_1\eta_2 + \eta_2\eta_3 + \eta_3\eta_1] + \frac{c_{44}}{2} [\eta_4^2 + \eta_5^2 + \eta_6^2] \\
& + \frac{\lambda_1}{2} (L_{1x}^2\eta_1 + L_{2y}^2\eta_2 + L_{3z}^2\eta_3) + \frac{\lambda'_1}{2} (L_{1x}^2 + L_{2y}^2 + L_{3z}^2) (\eta_1 + \eta_2 + \eta_3) \\
& + \frac{\lambda_2}{2} (L_{1y}L_{1z}\eta_4 + L_{2x}L_{2z}\eta_5 + L_{3x}L_{3y}\eta_6) + \frac{\lambda_3}{2} (M_x^2\eta_1 + M_y^2\eta_2 + M_z^2\eta_3) \\
& + \frac{\lambda'_3}{2} (M_yM_z\eta_4 + M_xM_z\eta_5 + M_xM_y\eta_6) \\
& + \frac{\lambda}{2} ([L_{1x}L_{1y} + L_{2x}L_{2y}] \eta_6 + [L_{1x}L_{1z} + L_{3x}L_{3z}] \eta_5 + [L_{2y}L_{2z} + L_{3y}L_{3z}] \eta_4) \\
& + \frac{A}{2} \left[\left(\frac{\partial \mathbf{L}_1}{\partial x} \right)^2 + \left(\frac{\partial \mathbf{L}_1}{\partial y} \right)^2 + \left(\frac{\partial \mathbf{L}_1}{\partial z} \right)^2 + \left(\frac{\partial \mathbf{L}_2}{\partial x} \right)^2 + \left(\frac{\partial \mathbf{L}_2}{\partial y} \right)^2 + \left(\frac{\partial \mathbf{L}_2}{\partial z} \right)^2 \right. \\
& \left. + \left(\frac{\partial \mathbf{L}_3}{\partial x} \right)^2 + \left(\frac{\partial \mathbf{L}_3}{\partial y} \right)^2 + \left(\frac{\partial \mathbf{L}_3}{\partial z} \right)^2 \right] \\
& + \frac{B}{2} \left[\left(\frac{\partial \mathbf{L}_1}{\partial x} \right)^2 + \left(\frac{\partial \mathbf{L}_2}{\partial y} \right)^2 + \left(\frac{\partial \mathbf{L}_3}{\partial z} \right)^2 \right] \\
& - \mathbf{M} \cdot \mathbf{H}.
\end{aligned} \tag{1}$$

Here, x , y and z are oriented along the axes of the cubic cell of UO_2 , and η_i , $i = 1, 6$ are the corresponding strain using Voigt notation.